

From Parametricity to Conservation Laws, via Noether's Theorem

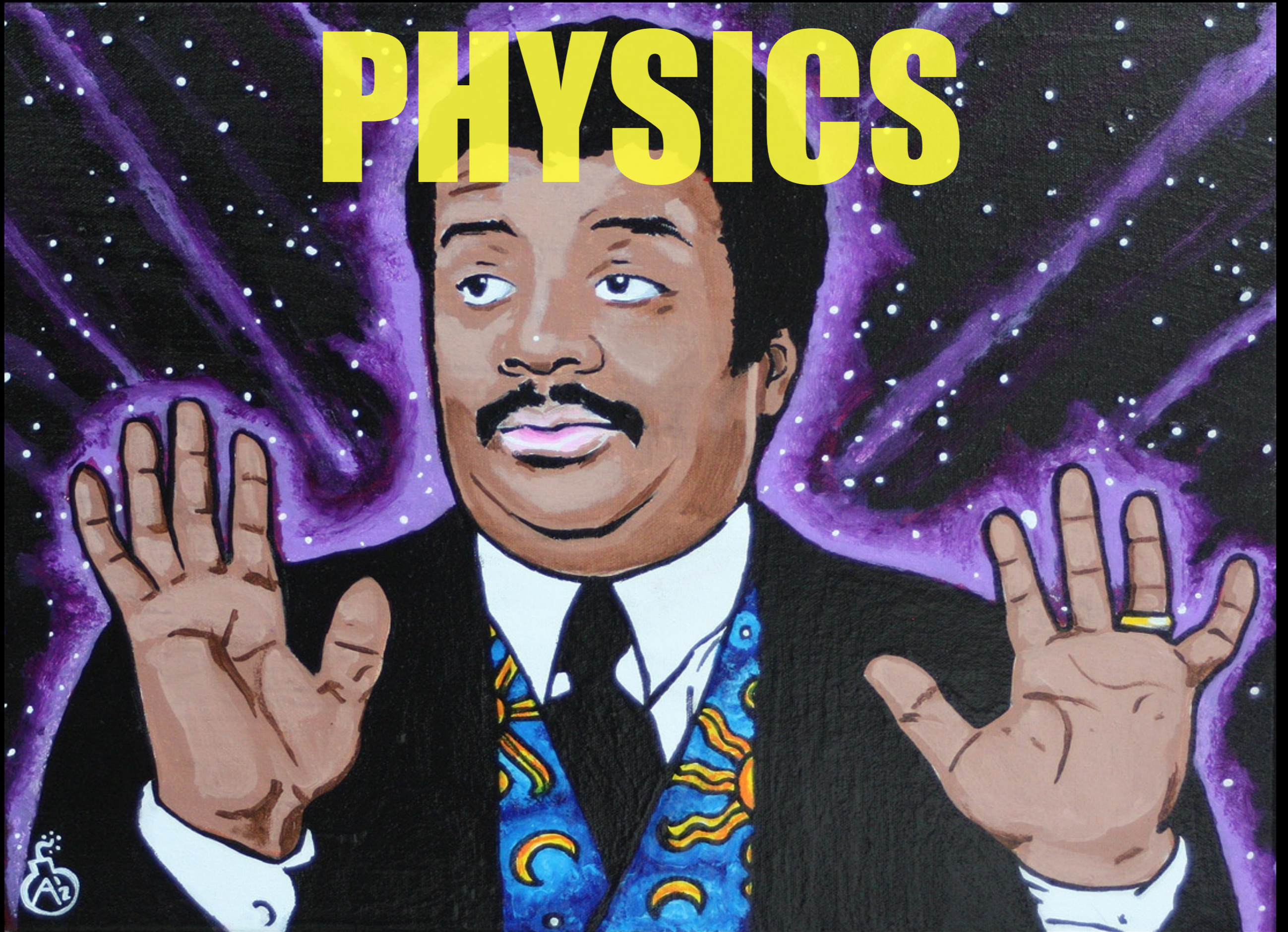
Written by Robert Atkey

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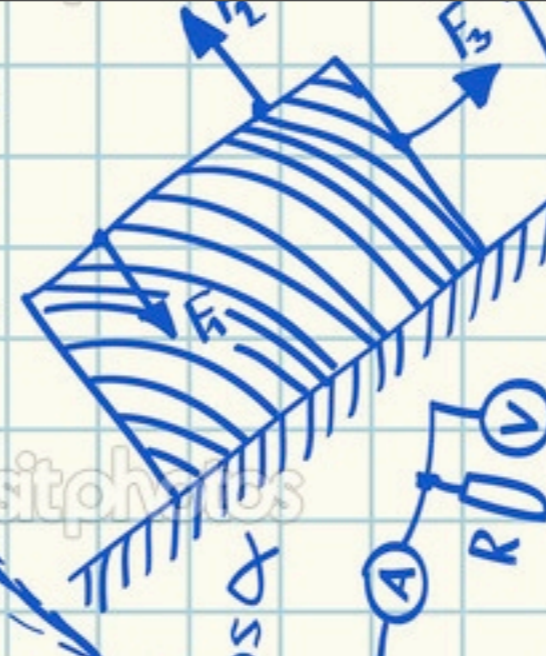
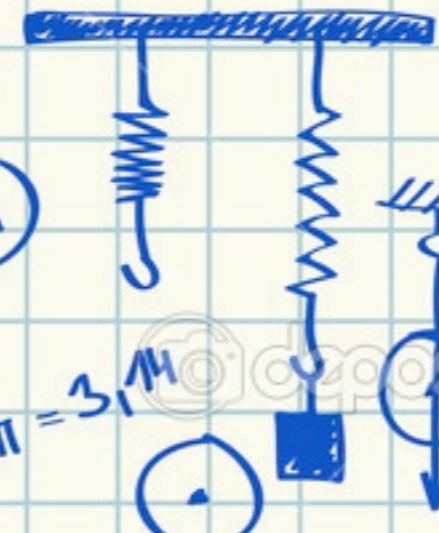
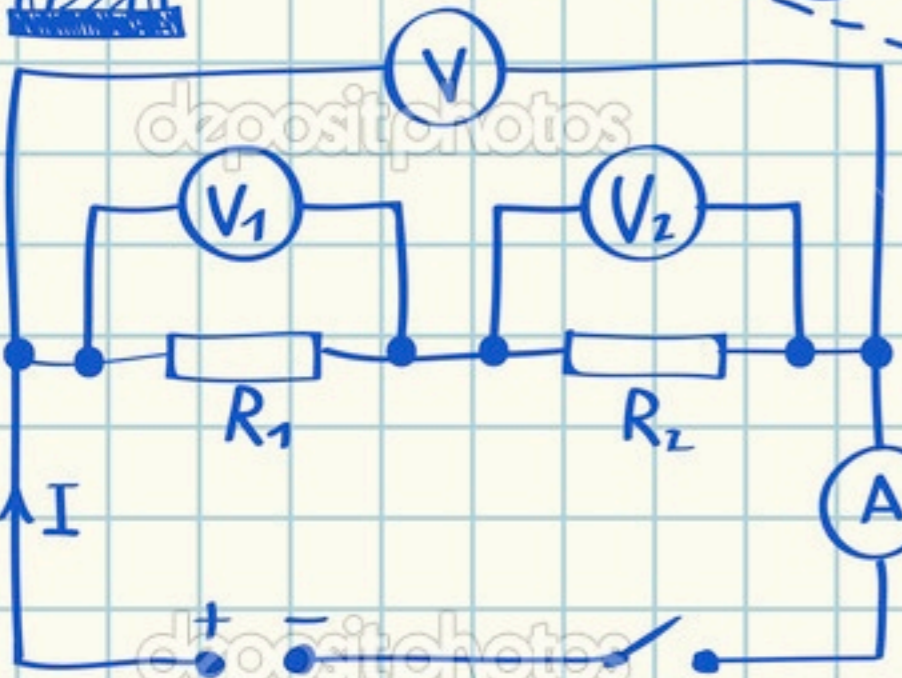
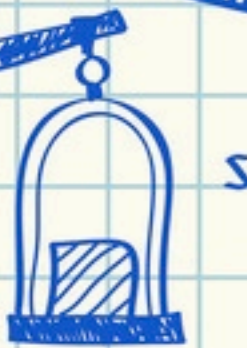
Presented by Ben Carriel & Ben Greenman

2014-04-07

PHYSICS

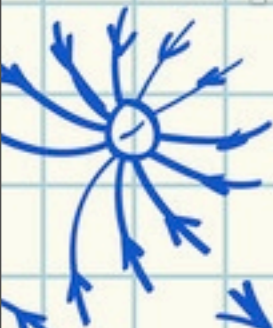


$$s = \frac{1}{2} at^2$$

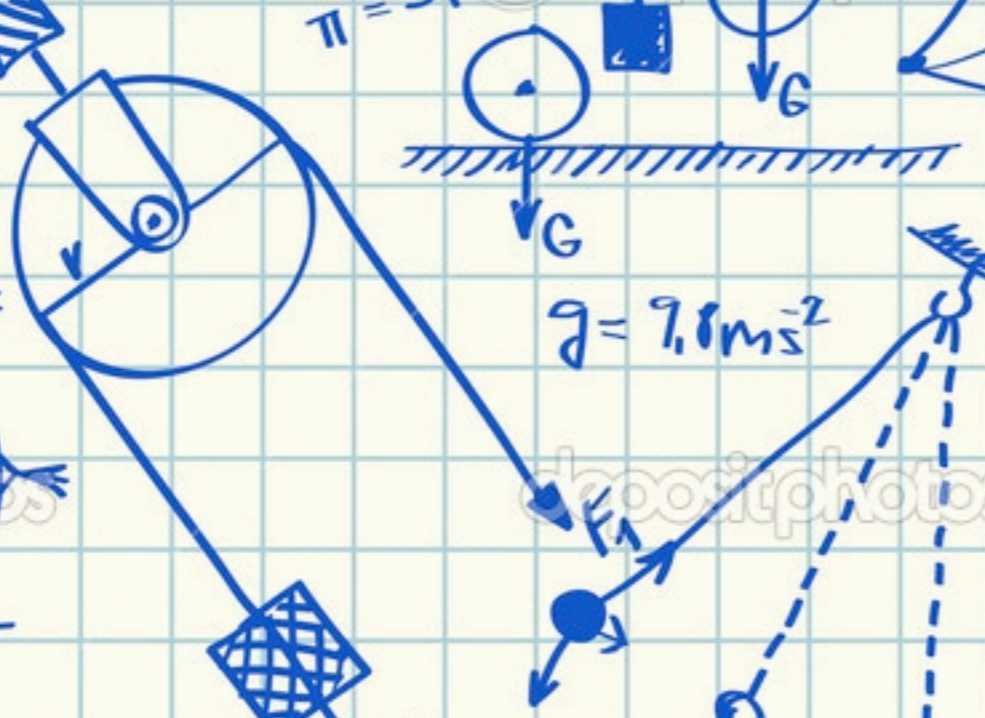
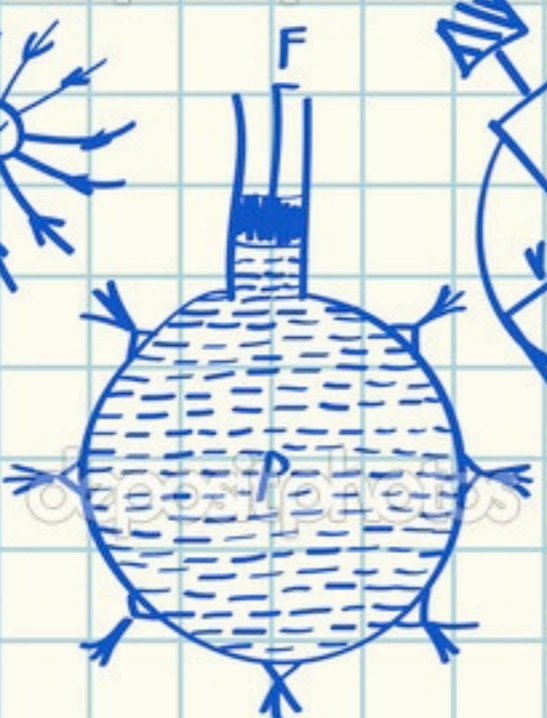


$$\Phi = BS \cos \alpha$$

$$F = ma$$

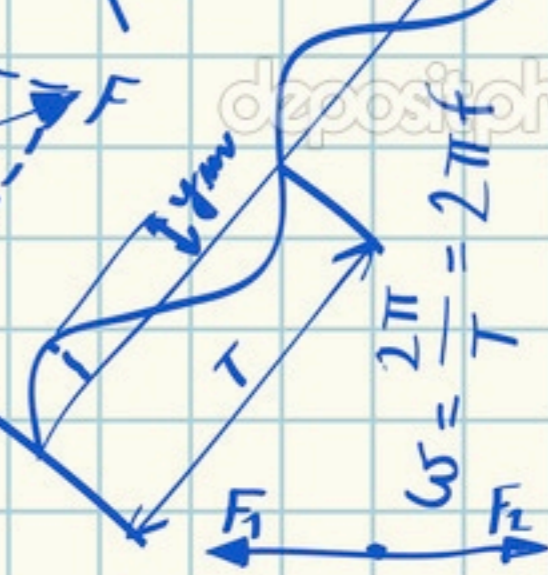


$$p = \rho h g$$

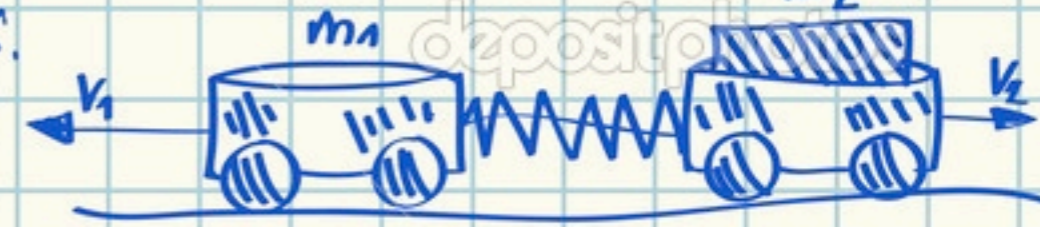


$$g = 9.8 \text{ m/s}^2$$

$$F = m v \sin \omega = n$$



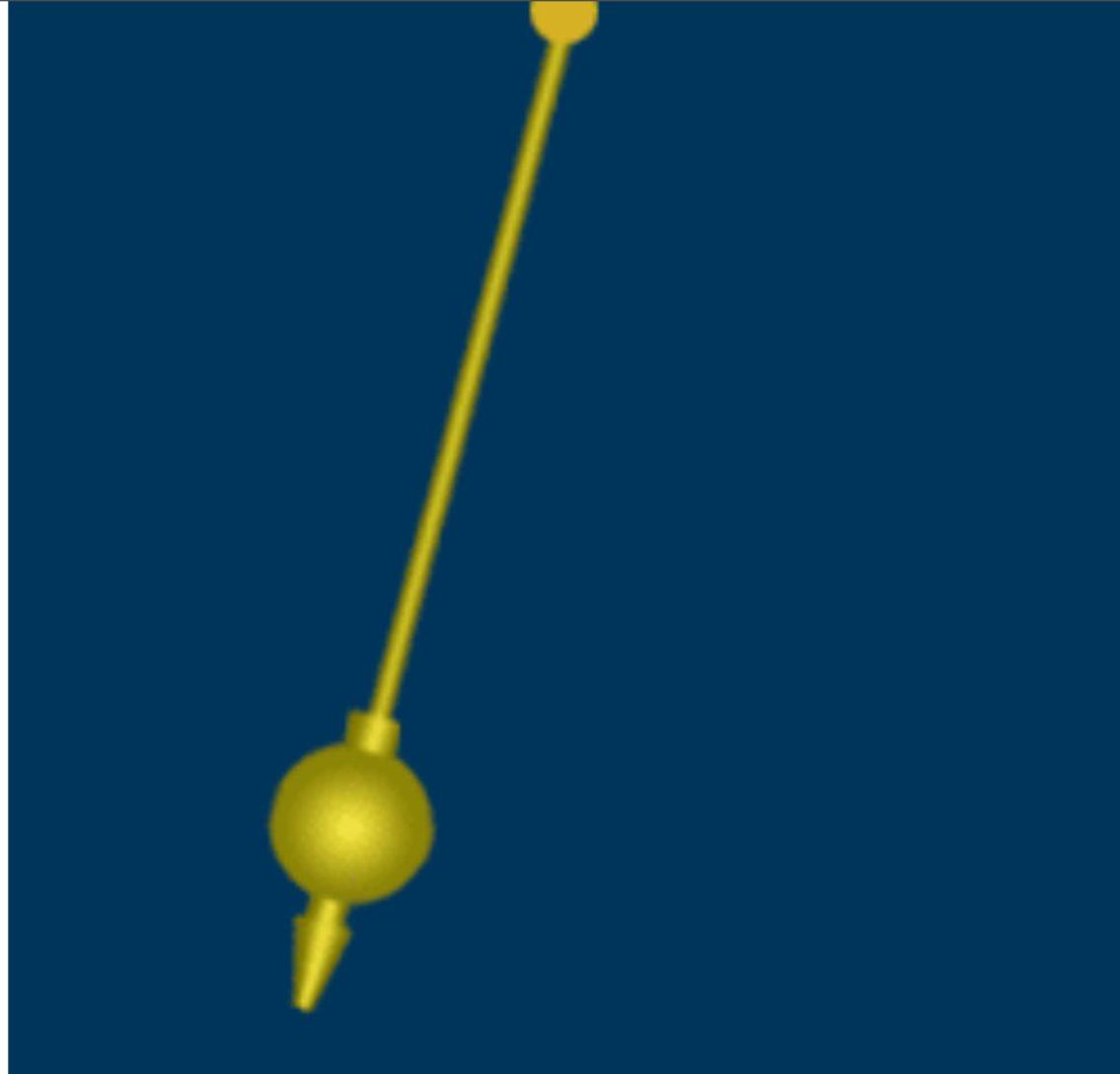
$$v = \frac{2\pi}{T}$$

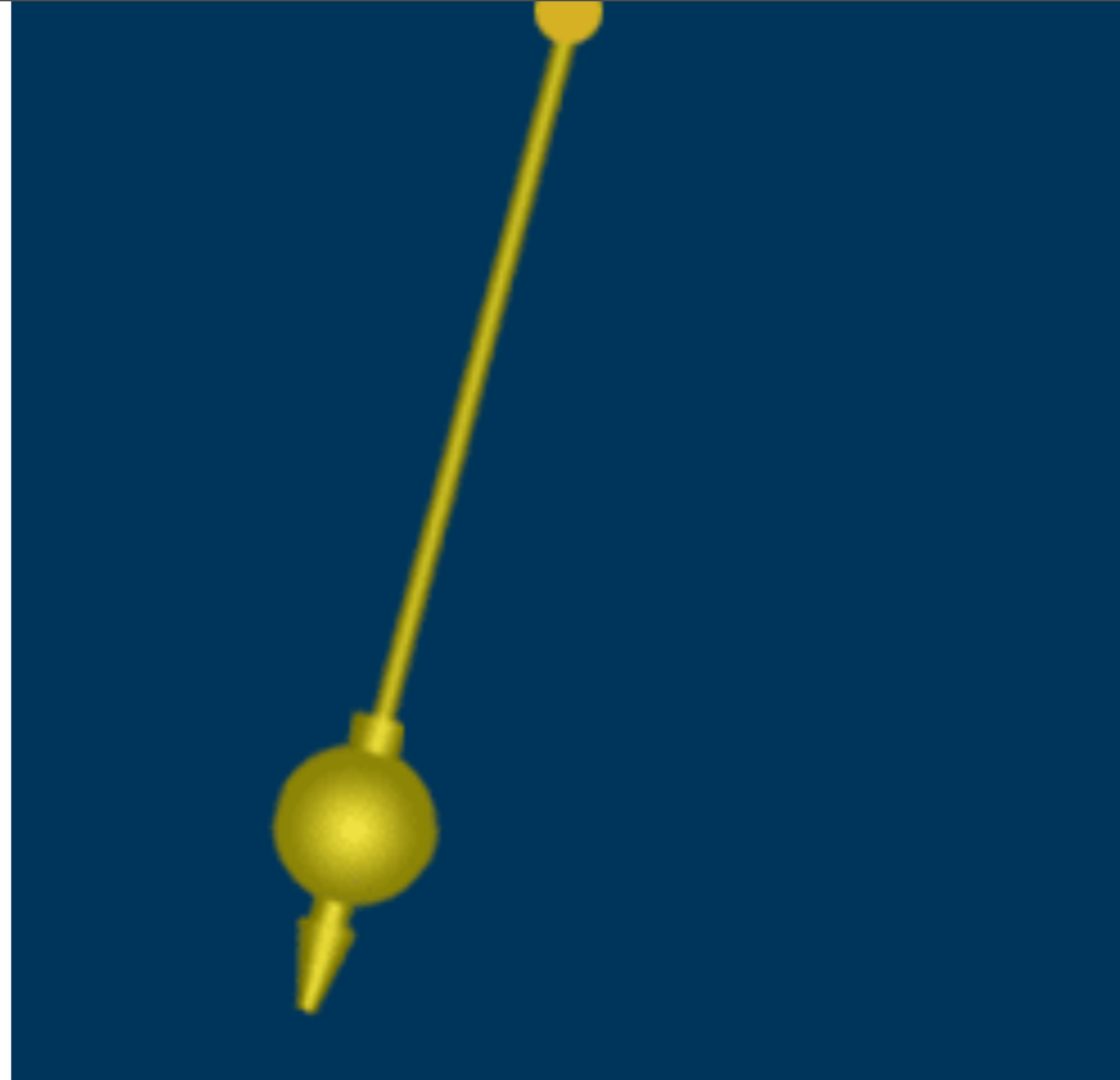


$$F = \alpha \frac{m_1 m_2}{r^2}$$

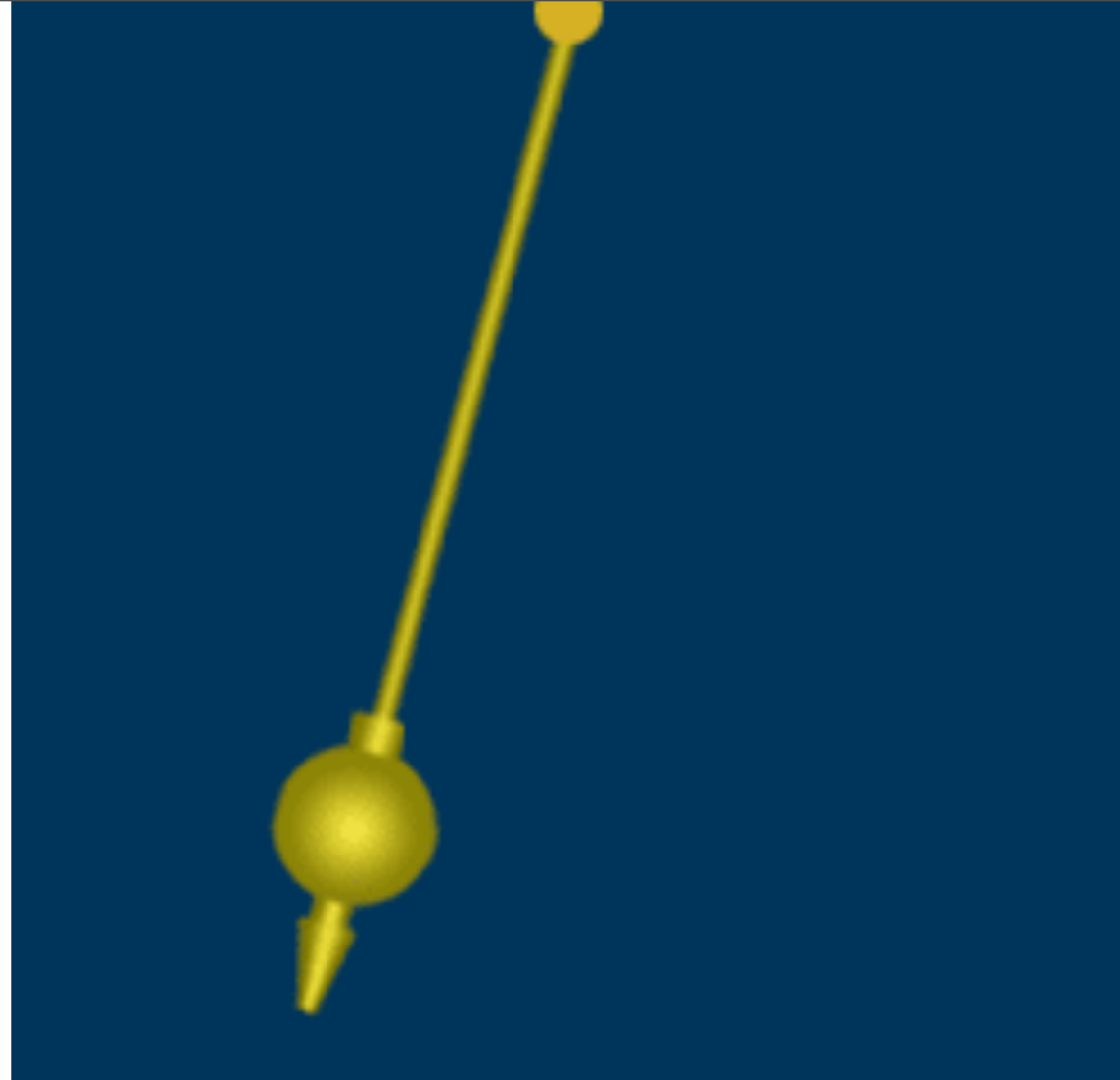
$$v = \sqrt{2\epsilon_0}$$







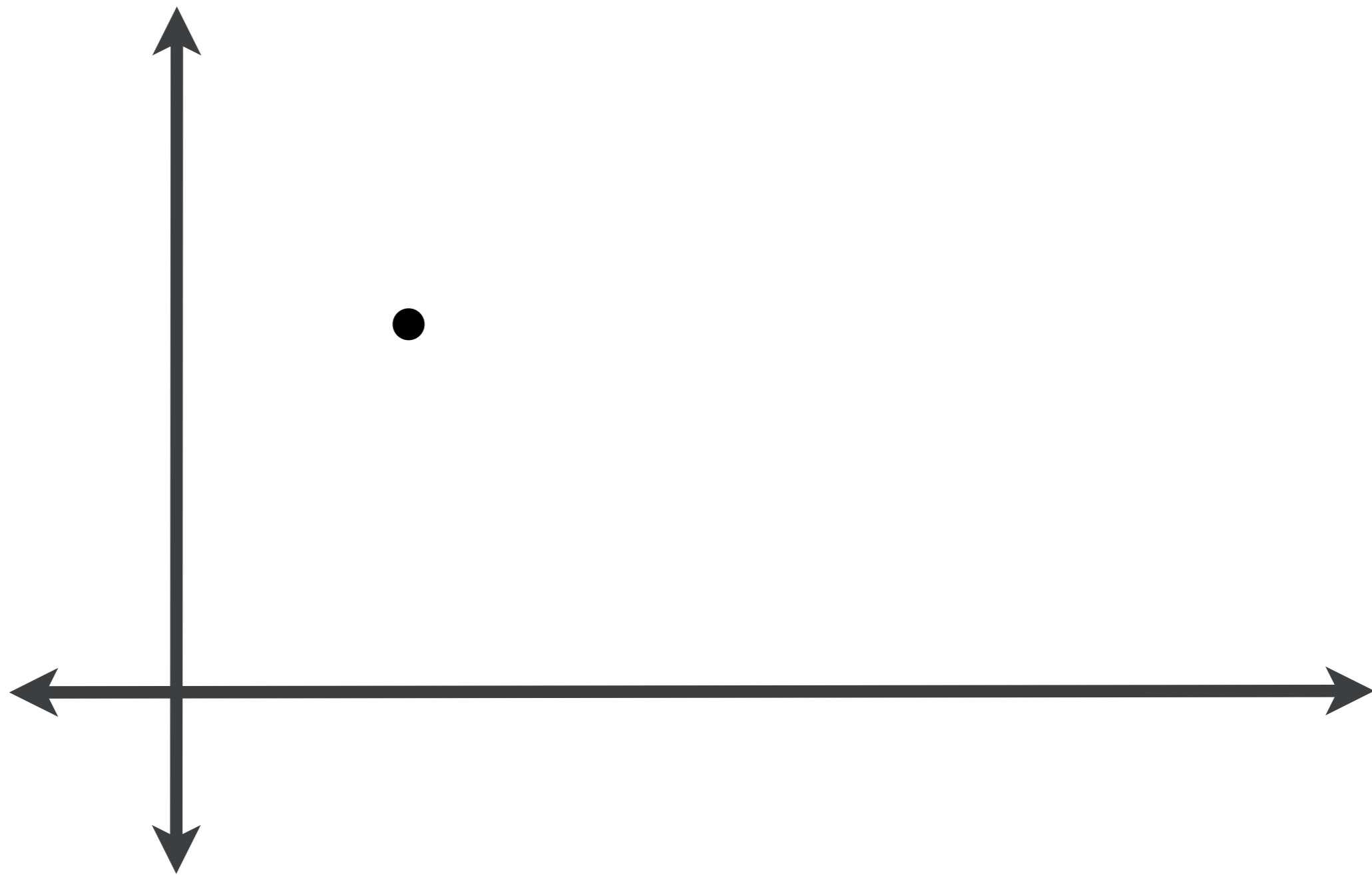
$$TE = mgh + \frac{1}{2}mv^2$$

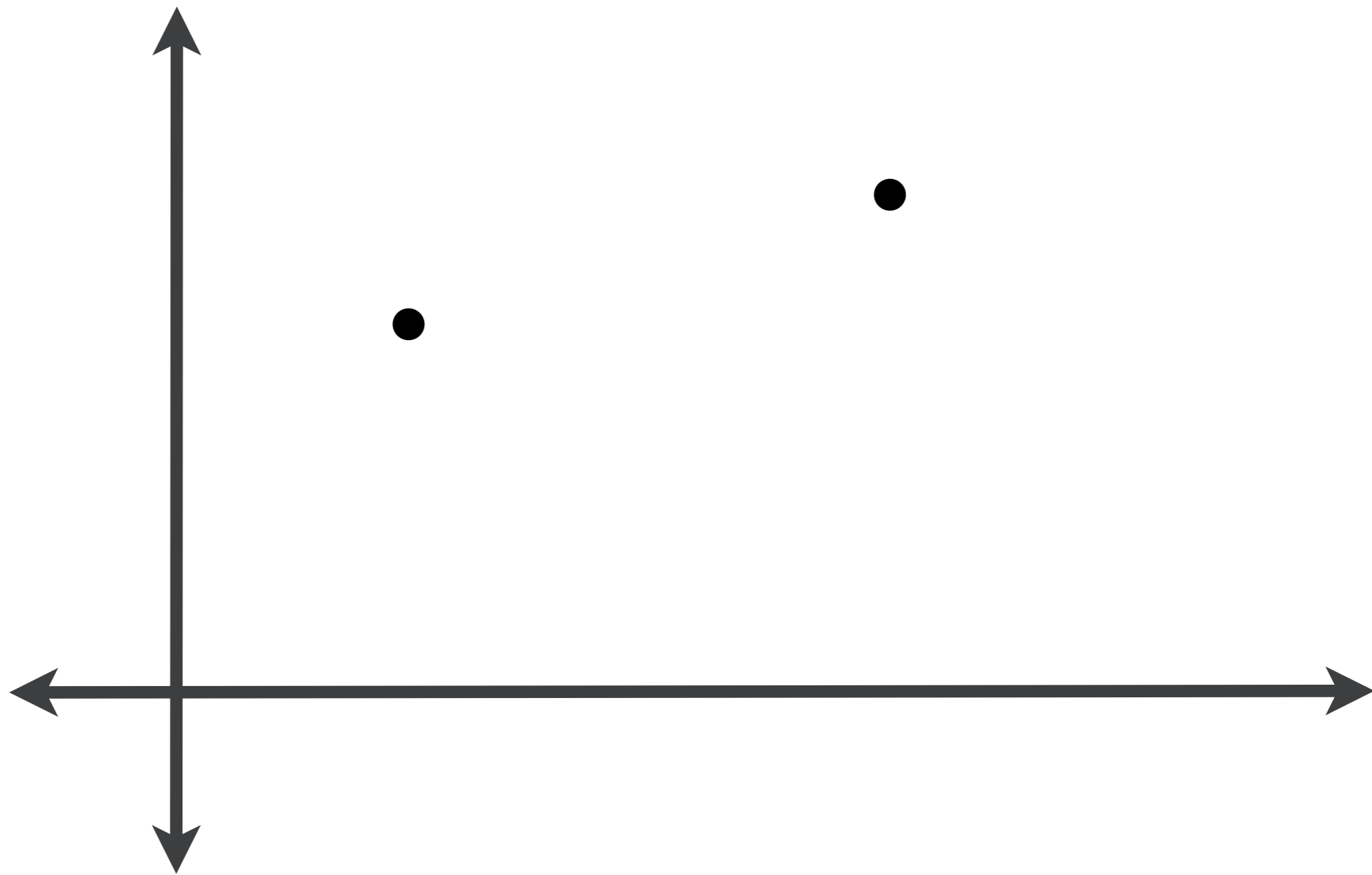


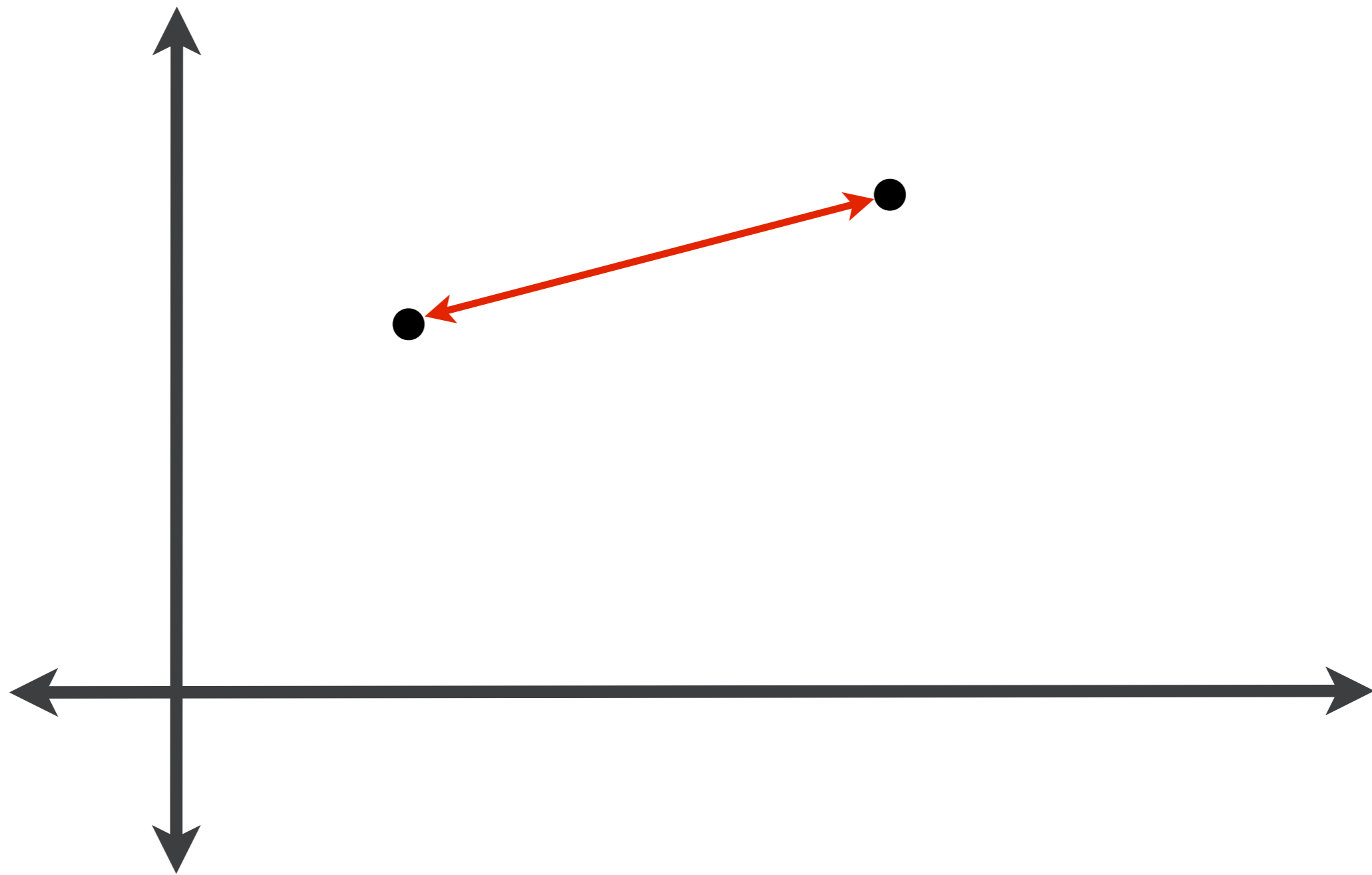
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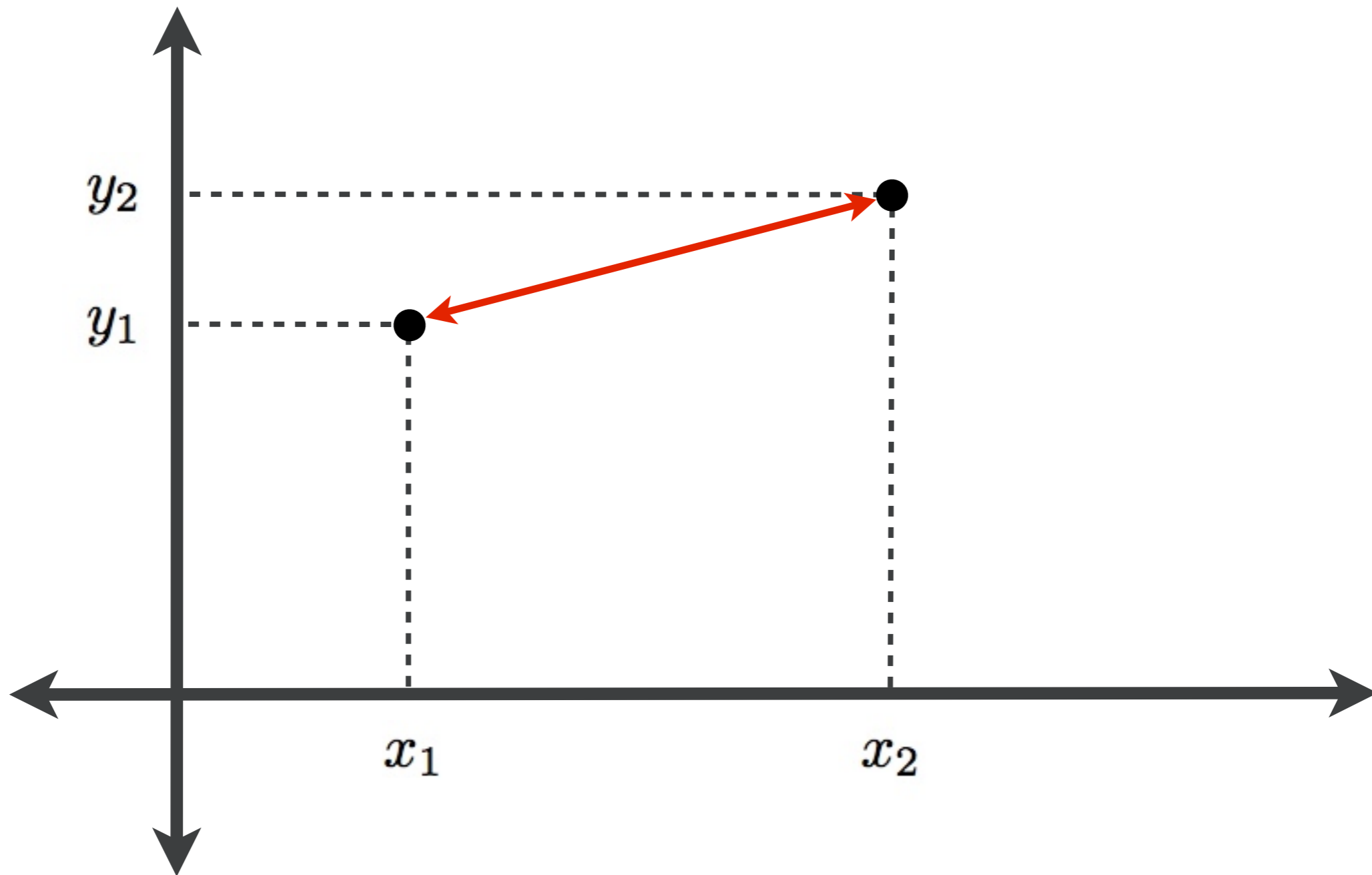
$$TE = PE + KE$$



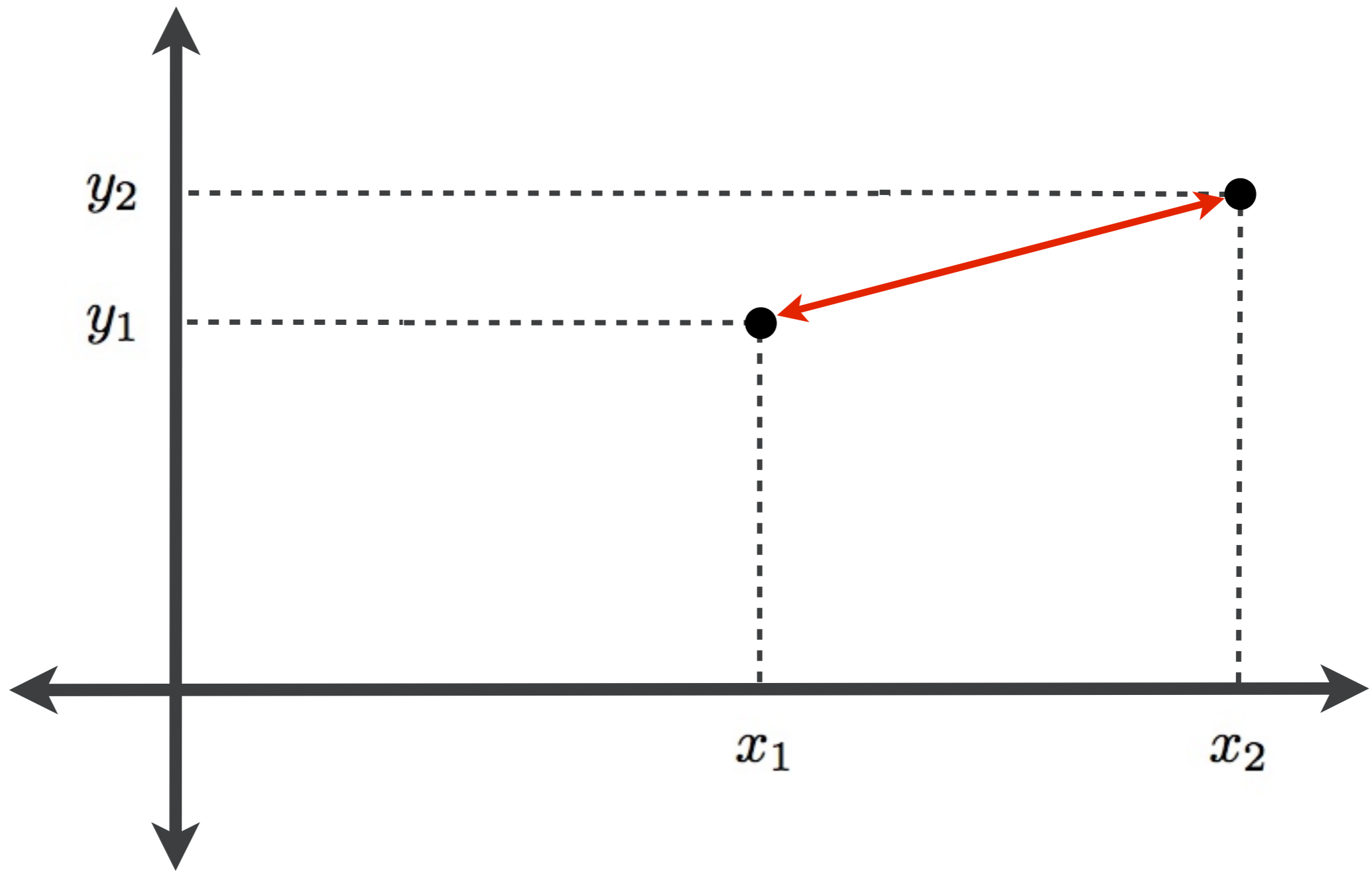






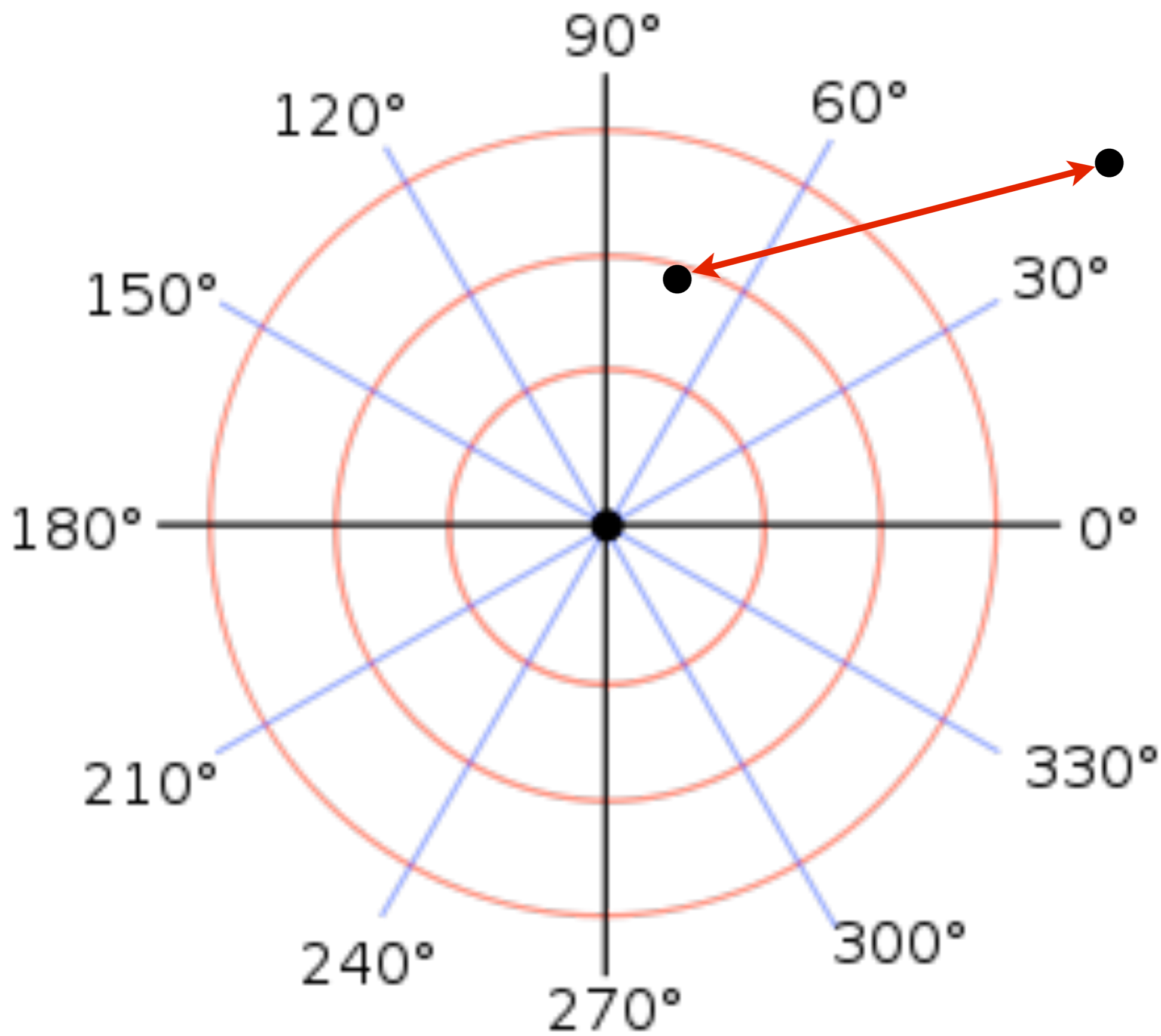


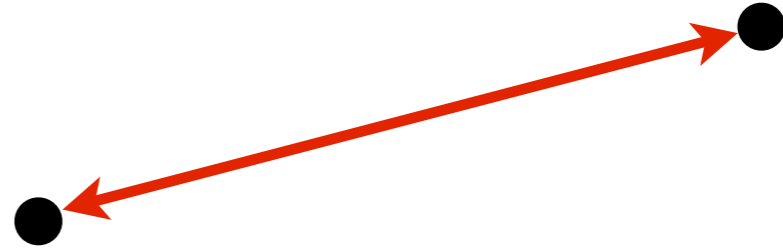
$$\text{distance} = \sqrt{(x_1 - x_2)^2 + (y_1 - y_2)^2}$$



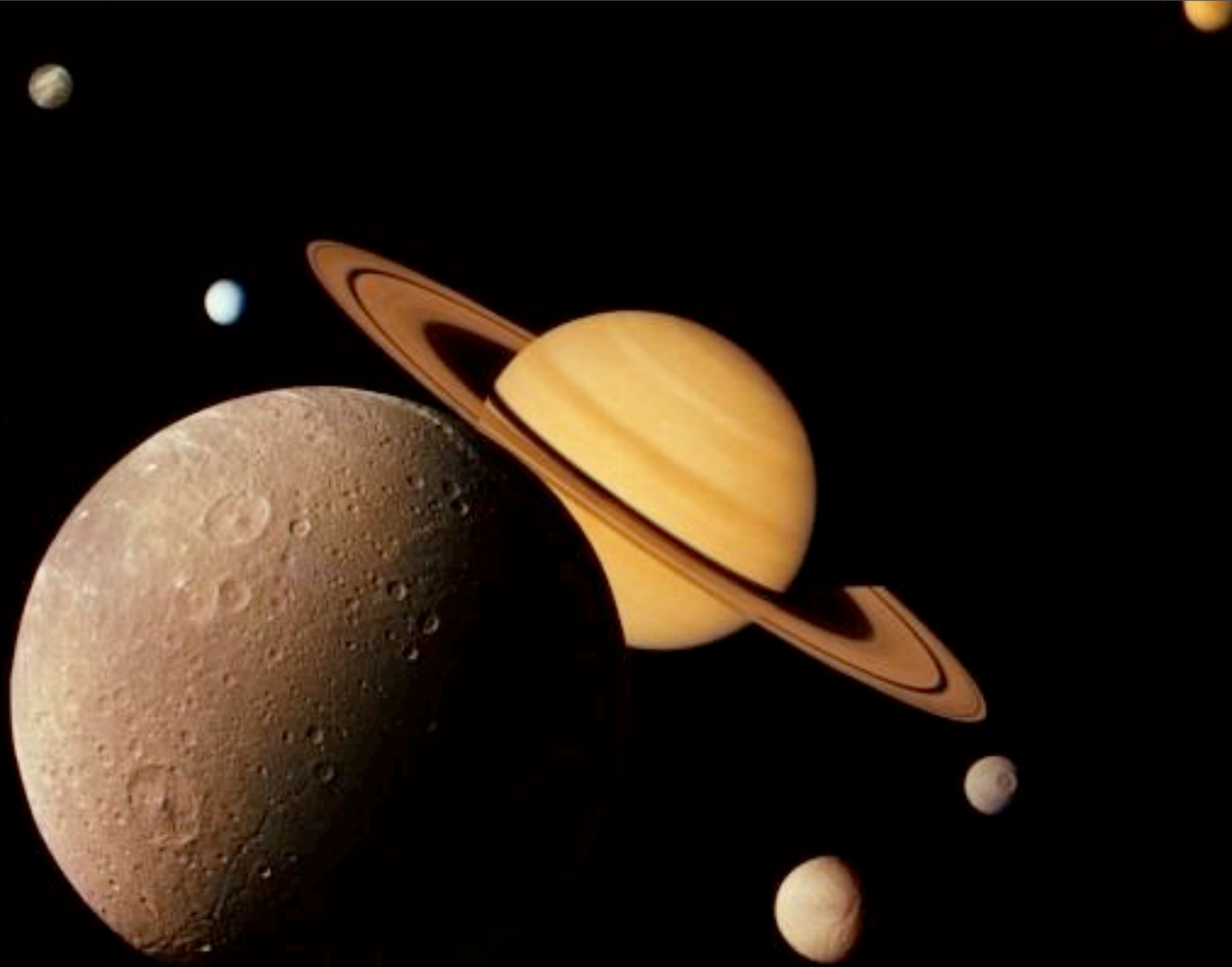
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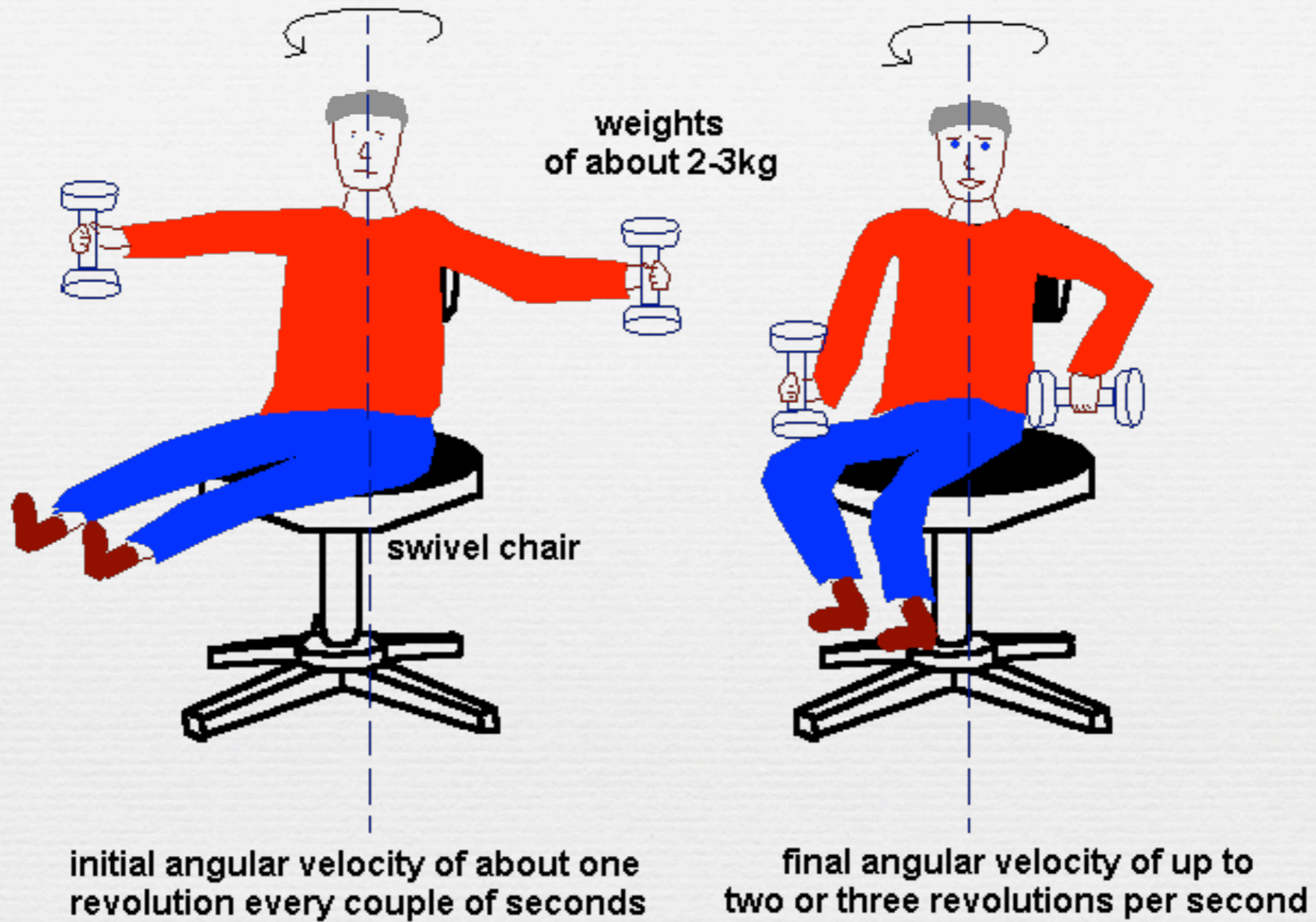






Distance is invariant under translation and change of coordinate representation.







Noether's Theorem

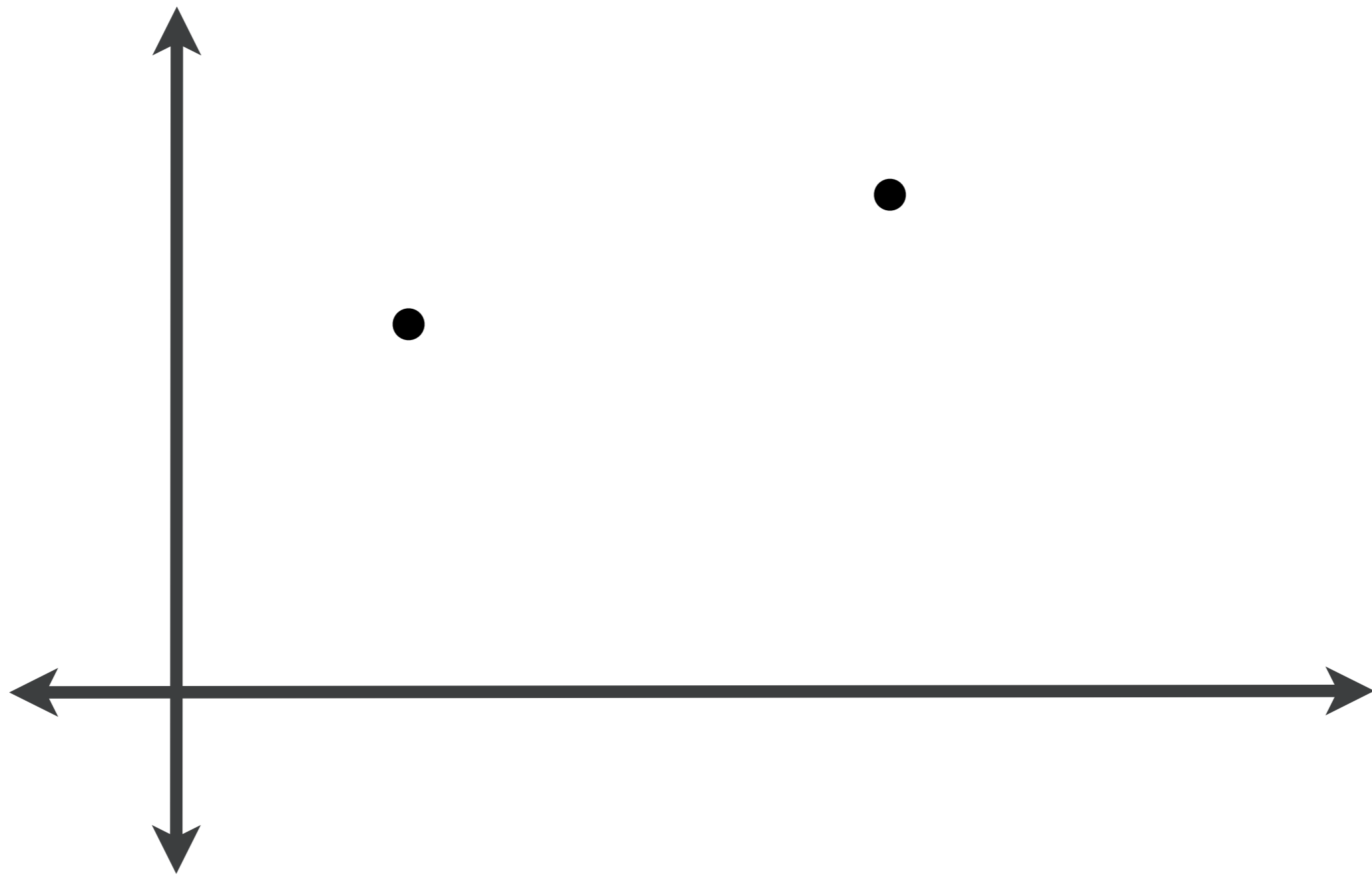


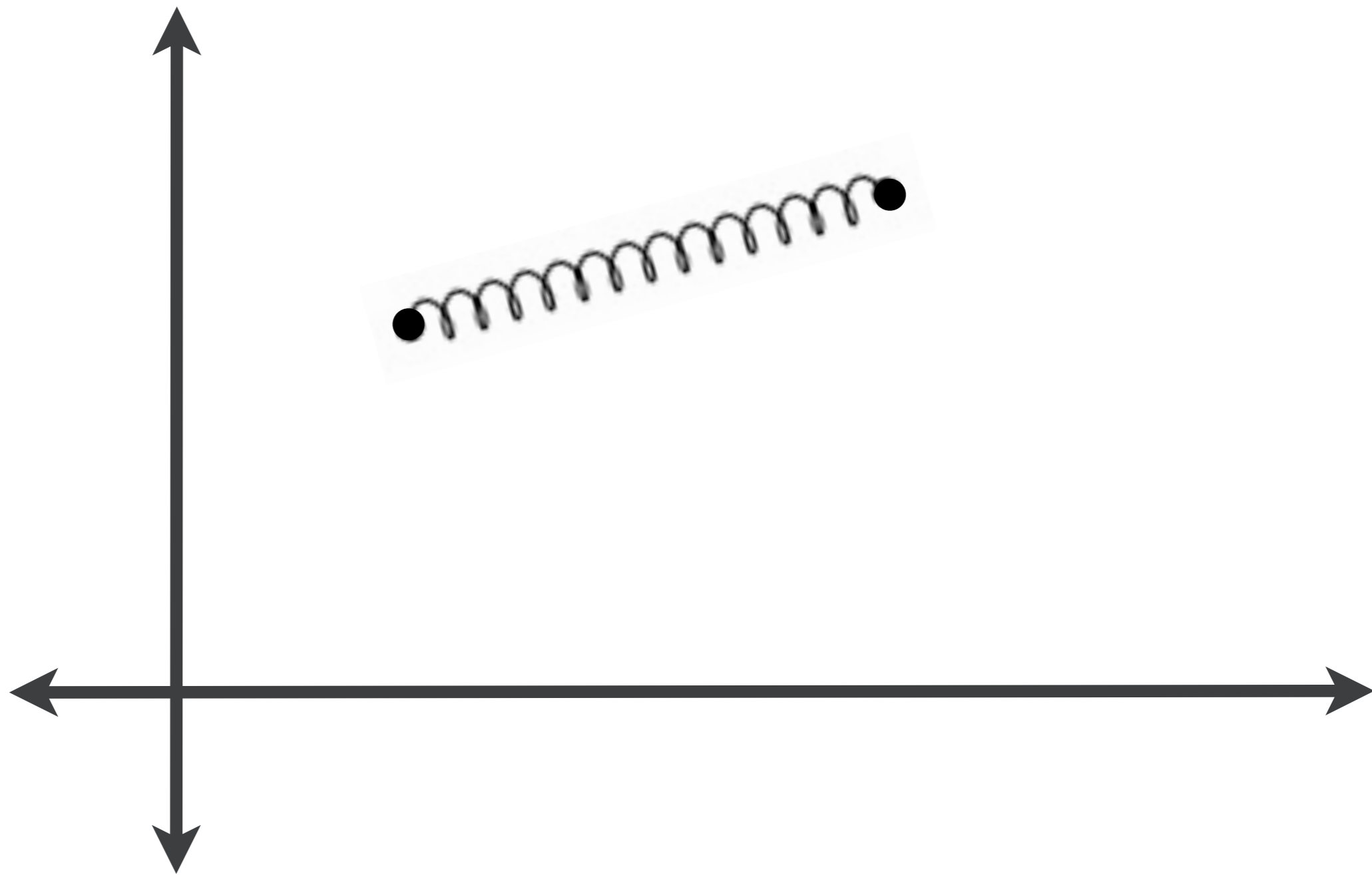
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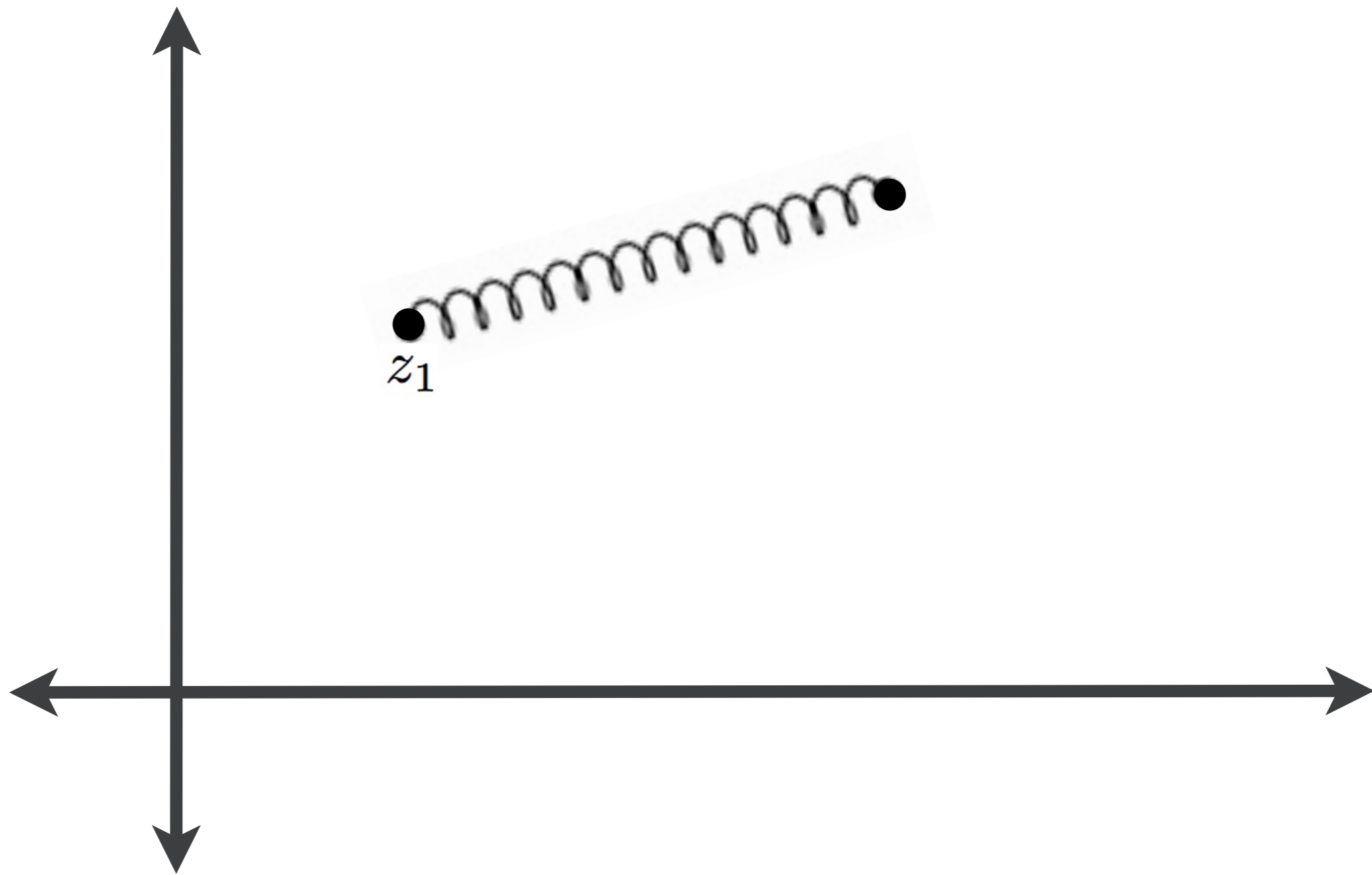
- (1915) "Any differentiable symmetry of the action of a physical system has a corresponding conservation law"

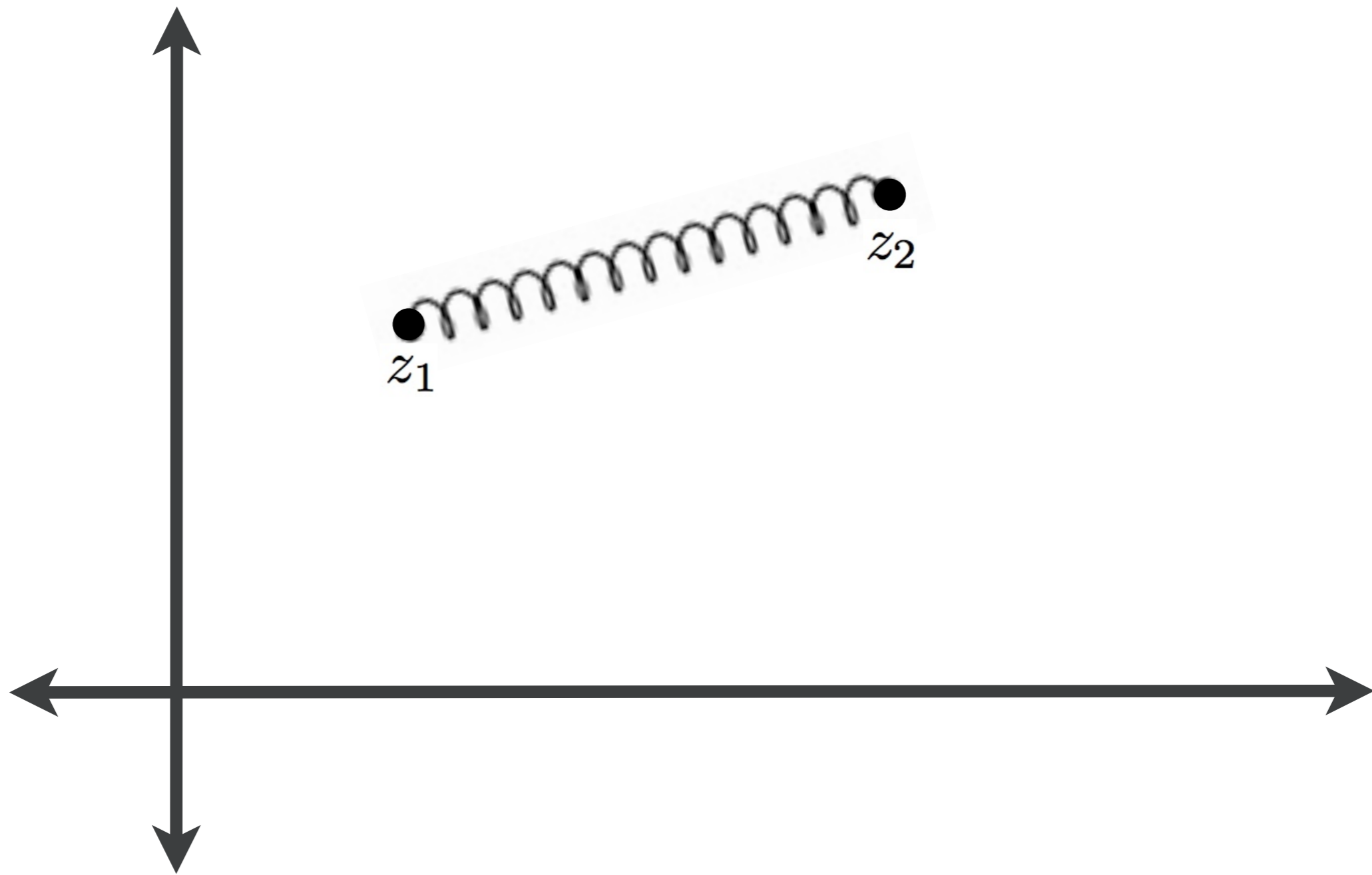


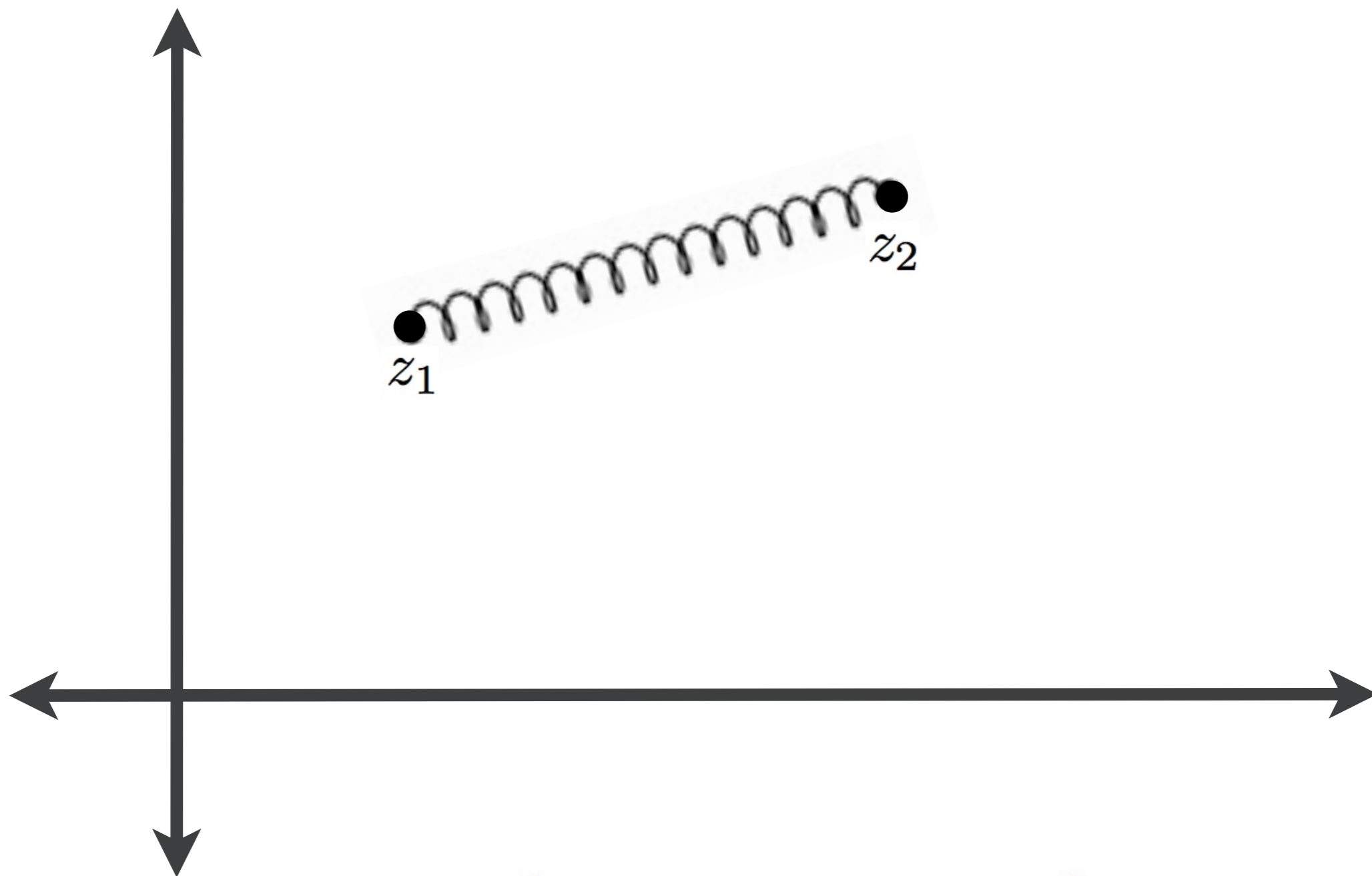
Quick Example











$$L(t, z_1, z_2, \dot{z}_1, \dot{z}_2) = \frac{1}{2}m(\dot{z}_1^2 + \dot{z}_2^2) - \frac{1}{2}k(z_1 - z_2)^2$$

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$$\forall d \in \mathbb{R}^2$$

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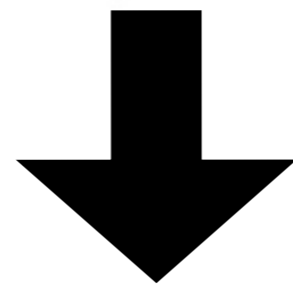
$$L(t, z_1, z_2, \dot{z}_1, \dot{z}_2) = \frac{1}{2}m(\dot{z}_1^2 + \dot{z}_2^2) - \frac{1}{2}k(z_1 - z_2)^2$$

$$\forall d \in \mathbb{R}^2 \quad L(t, z_1, z_2, \dot{z}_1, \dot{z}_2) = L(t, z_1 + d, z_2 + d, \dot{z}_1, \dot{z}_2)$$

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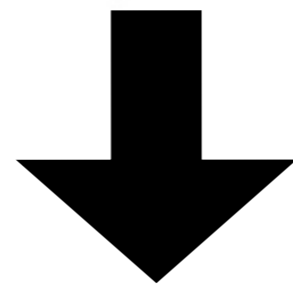
(Noether's Theorem)

$$\frac{d}{dt}m(\dot{x}_1 + \dot{x}_2) = 0$$

invariant under translation

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(Noether's Theorem)

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Conservation of Momentum

If the action

$$\mathcal{S}[q; a; b] = \int_a^b L(t, q, \dot{q}) dt$$

is invariant under Φ_ϵ and Ψ_ϵ , then

$$\frac{d}{dt} \left(\sum_{i=1}^n \frac{\partial L}{\partial \dot{q}_i} \psi_i + \left(L - \sum_{i=1}^n \dot{q}_i \frac{\partial L}{\partial \dot{q}_i} \right) \phi \right) = 0$$

where $\phi = \left. \frac{\partial \Phi}{\partial \epsilon} \right|_{\epsilon=0}$ and $\psi = \left. \frac{\partial \Psi}{\partial \epsilon} \right|_{\epsilon=0}$

Pretty cool, right?



TYPES

$(\lambda x : \text{unit} . 42 : \text{int})$

τ

$\pi \triangleq \Lambda T1 . \Lambda T2 . \Lambda T3 .$

$\lambda v : T1 \times T2 \times T3 .$

$v[T1] (\lambda x : T1 . \lambda y : T2 .$

$\lambda f : (\text{int} \rightarrow \text{int}) \text{ref} . \lambda n : \text{int} .$

$f := (\lambda \text{acc} : \text{int ref} . \lambda m : \text{int} .$

$\text{case } (n = m : \text{bool}) \text{ of}$

$(\text{acc} := (\text{mul } !\text{acc } m); \text{acc}) : \text{int}$

$(!f (\text{acc} := (\text{mul } !\text{acc } m); \text{acc}) (m+1)) : \text{int}$

$) (\text{ref } 1) 1) (\text{ref } \lambda x : \text{int} . x))$

Atkey (2014)

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- Define a type system for Lagrangian Mechanics.

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- Derive conservation laws as "free theorems" by parametricity.

Lagrangian: $L(t, z_1, z_2, \dot{z}_1, \dot{z}_2) = \frac{1}{2}m(\dot{z}_1^2 + \dot{z}_2^2) - \frac{1}{2}k(z_1 - z_2)^2$

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$\mathbb{R}^n \langle g, f \rangle$ \longrightarrow n -dimensional vectors of real numbers that vary with linear transformation g and translation f .

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Theorem (Noether). *Let $L(x, u, D_u^1, \dots, D_u^n)$, be a Lagrangian for $A \subseteq \mathbb{R}^n$, let $\varphi \in \text{Aut}(A)$ be a symmetry of A such that*

$$\varphi(L) + LD^i(\xi) = D^i(B^i) \quad B^i \in A$$

Then the Euler-Lagrange equations admit a conservation law $\forall i. D^i(C^i) = 0$.

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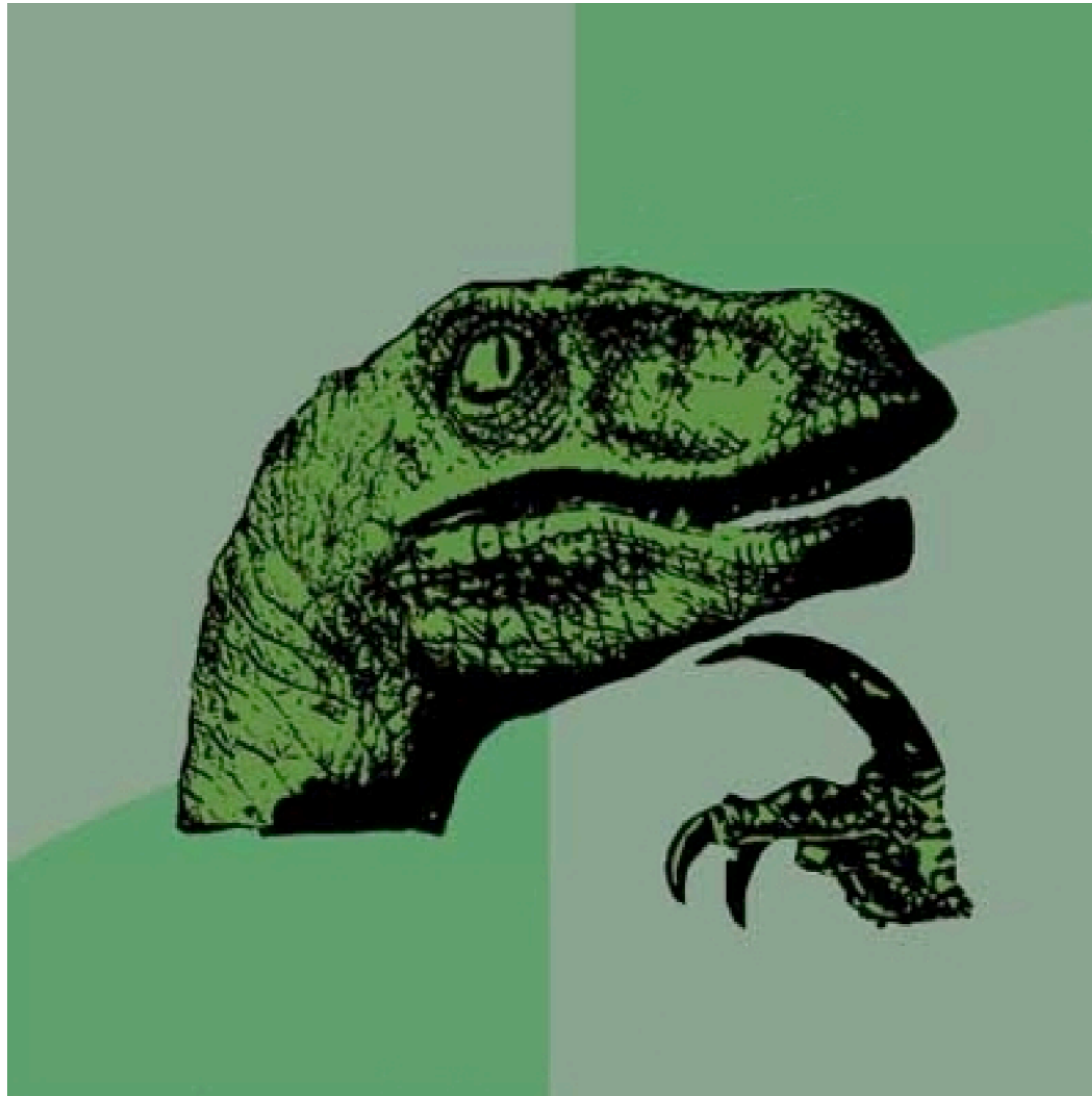
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Key point: we need an automorphism (i.e. symmetry) to start with

What does this mean?



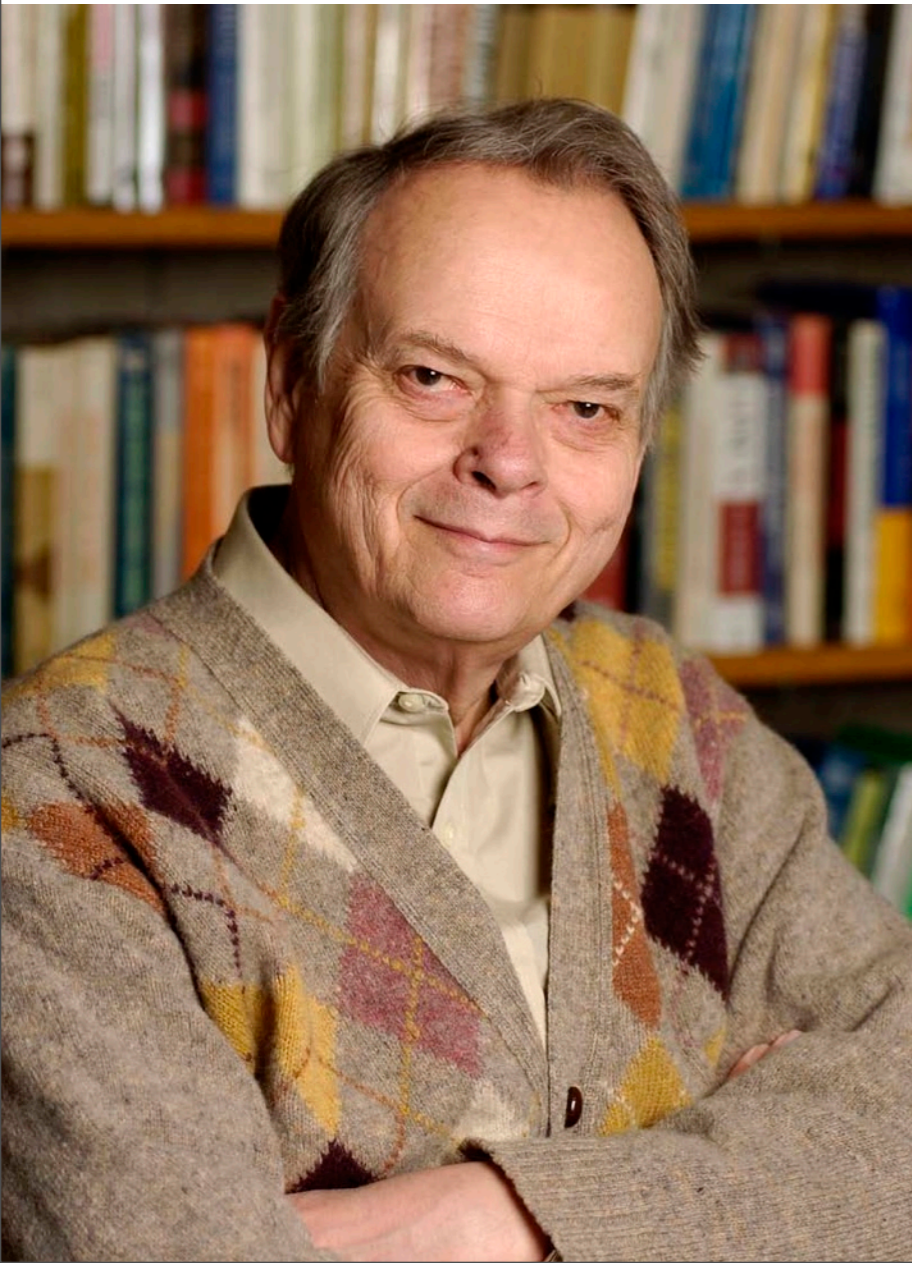
Reynolds:
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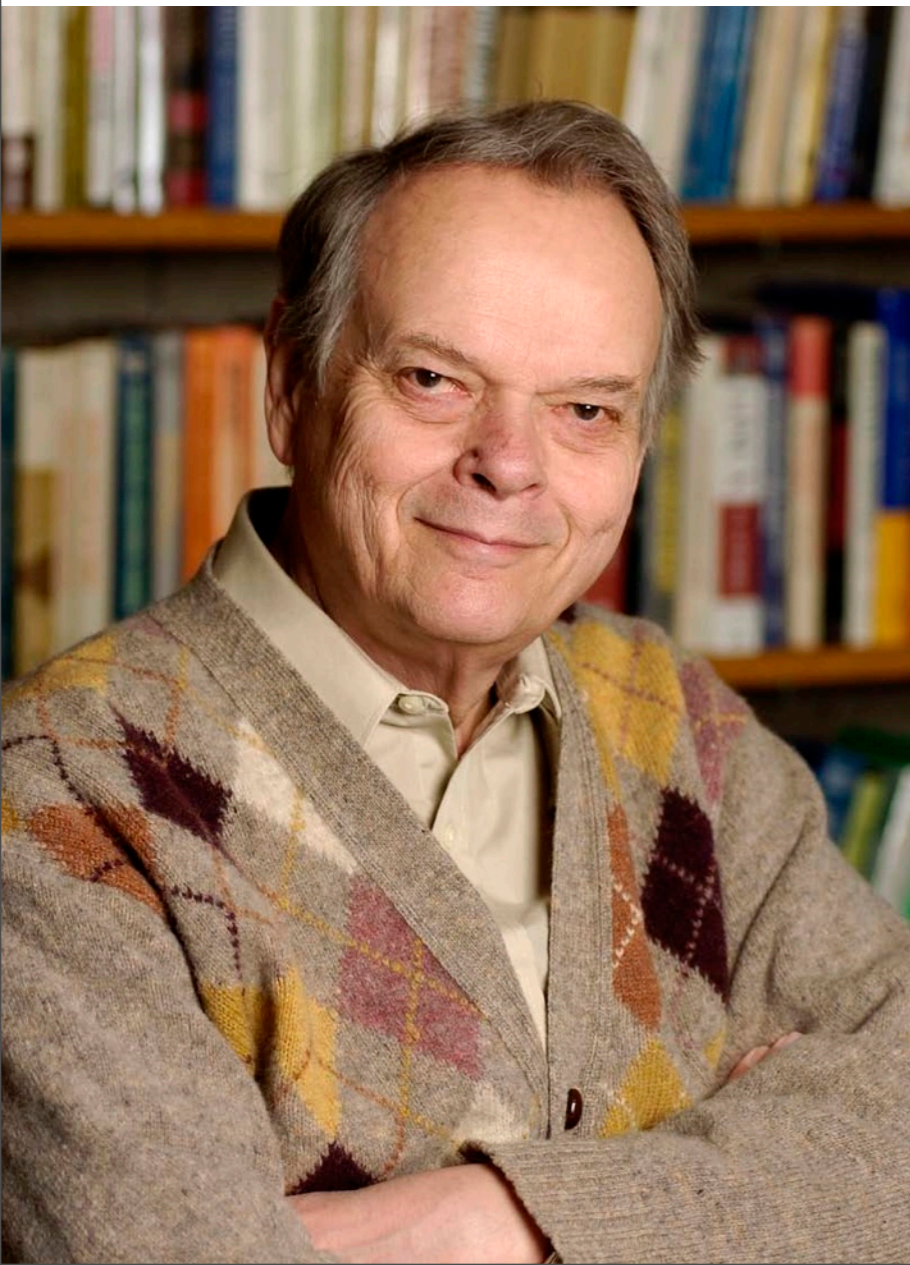
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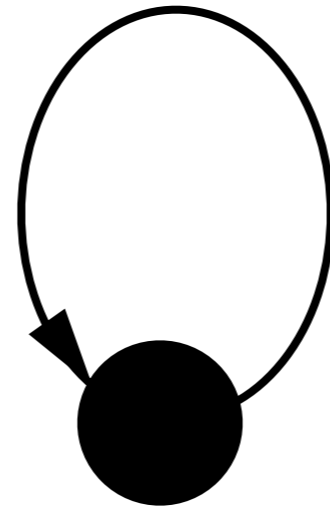
Atkey:
free theorems
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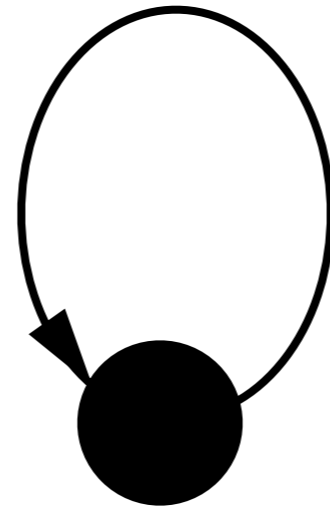


Atkey gives us a
geometric interpretation
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We'll argue: *Atkey* subsumes Reynolds + Wadler





Kinds are reflexive graphs

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- Reynolds: types are sets, parametricity comes from the relations between them.

Kinds are reflexive graphs

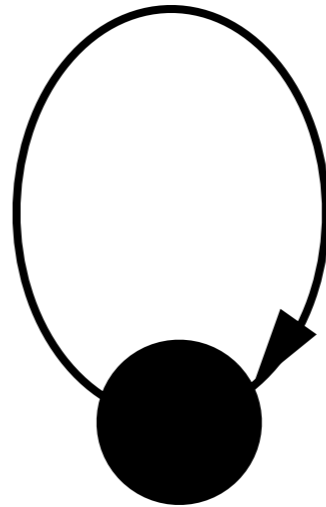
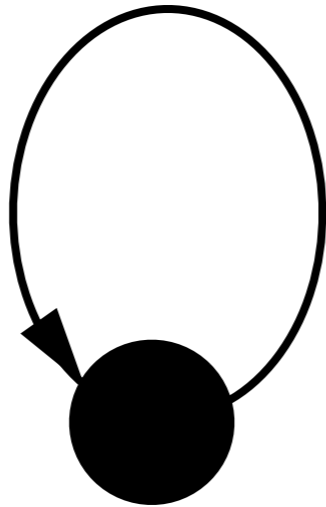
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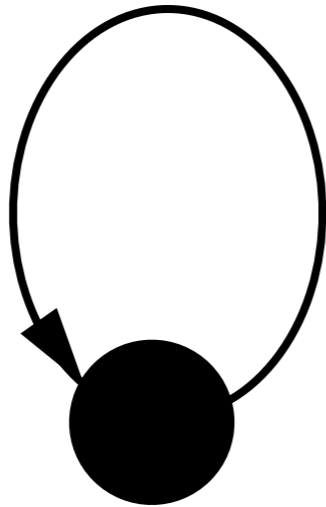
- Reynolds: types are sets, parametricity comes from the relations between them.
- Basic relation between Reynolds' types is the subset relation (\subseteq).
- Form a graph where the objects are types and the edges order types by \subseteq .

Example: bool

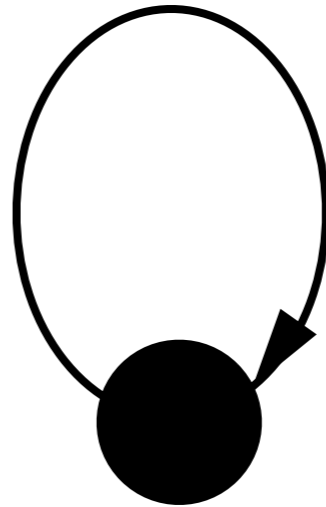
Example: bool



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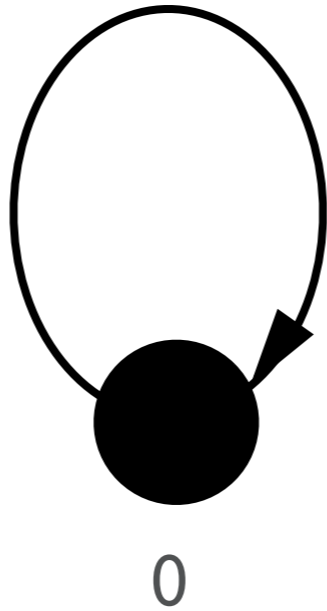
True



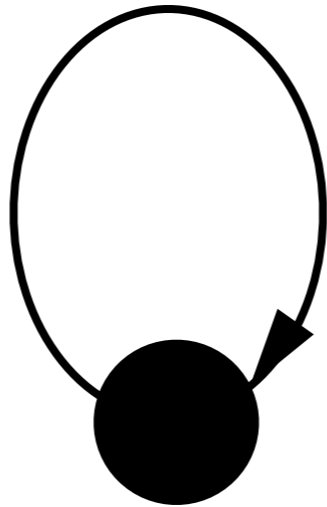
False

Example: nat

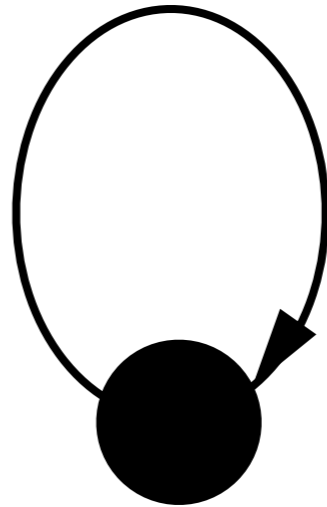
Example: nat



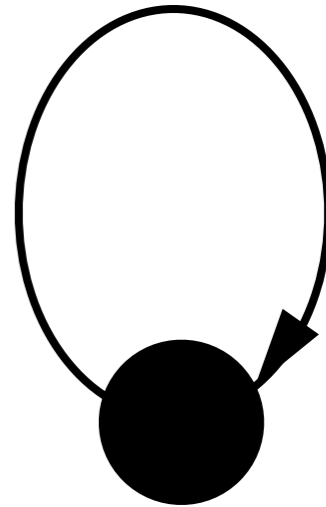
Example: nat



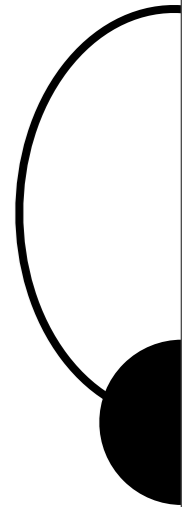
0



1



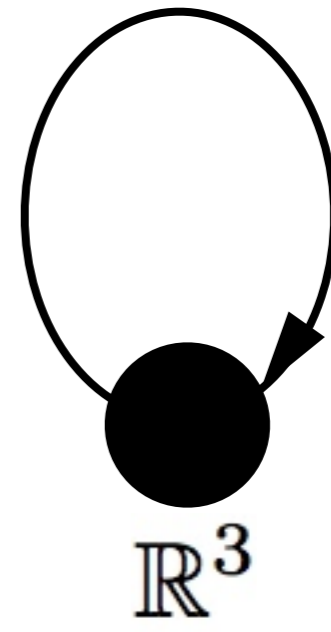
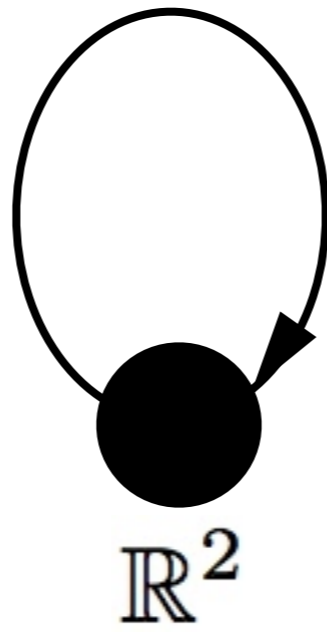
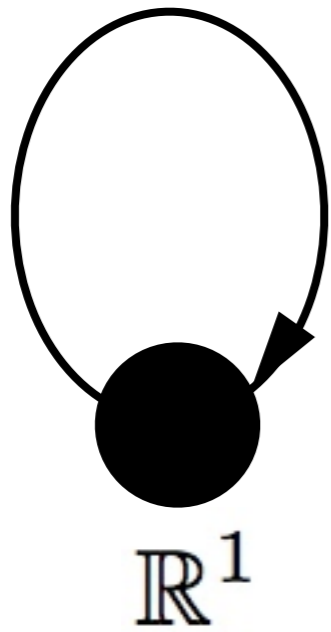
2



3

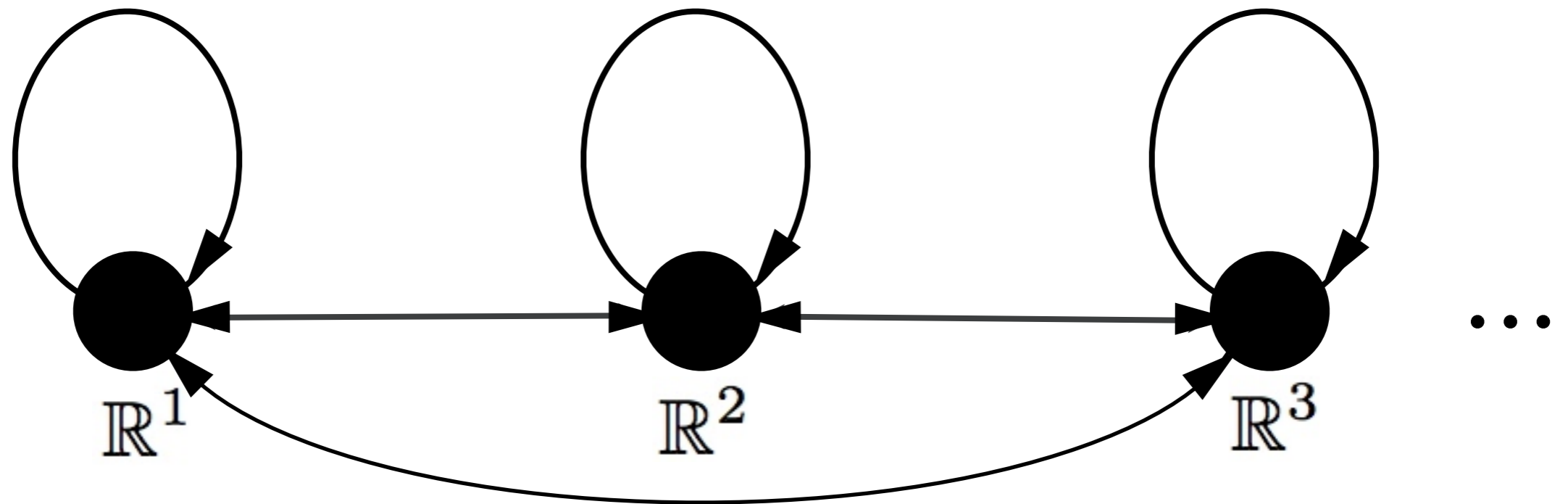
Example: cartesian space $(\mathbb{R}^1 \dots \mathbb{R}^n)$

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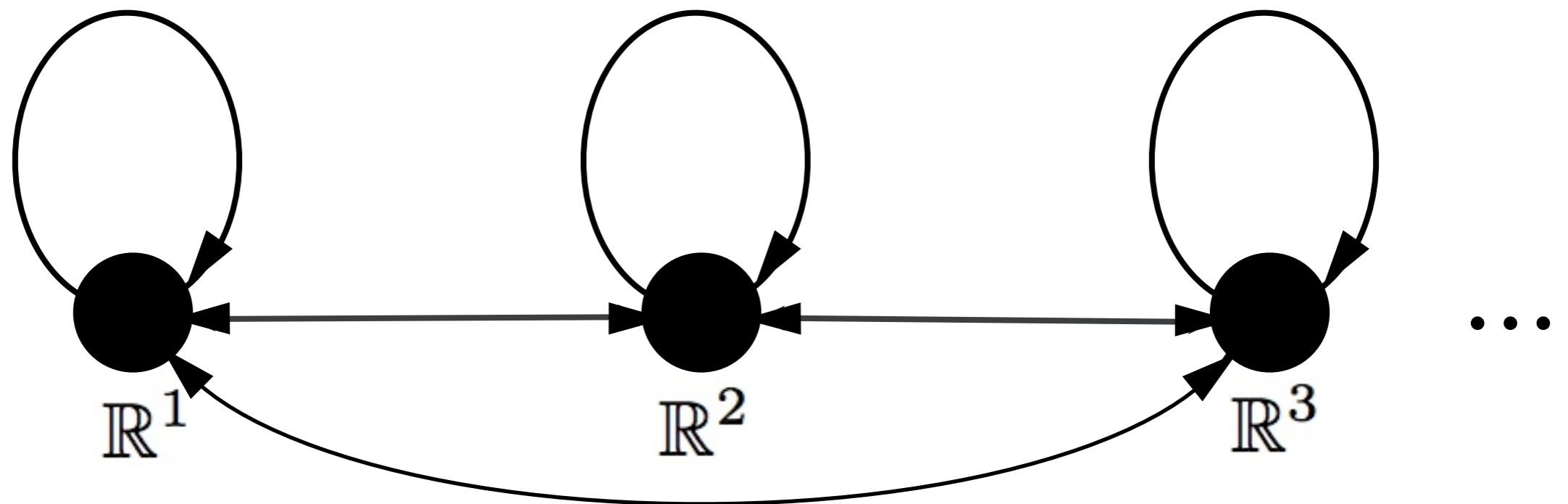


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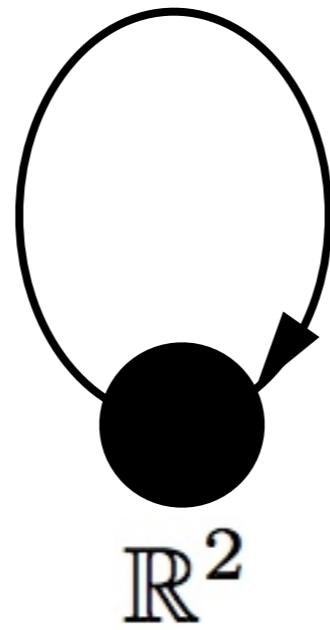


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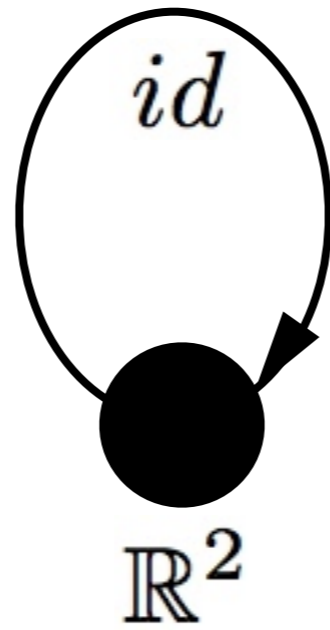


Each arrow represents a family of diffeomorphisms

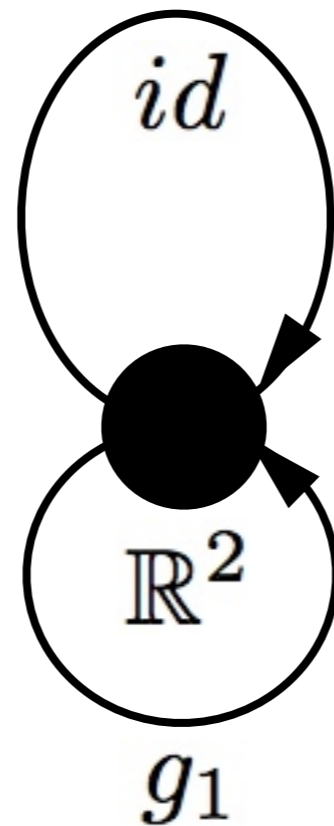
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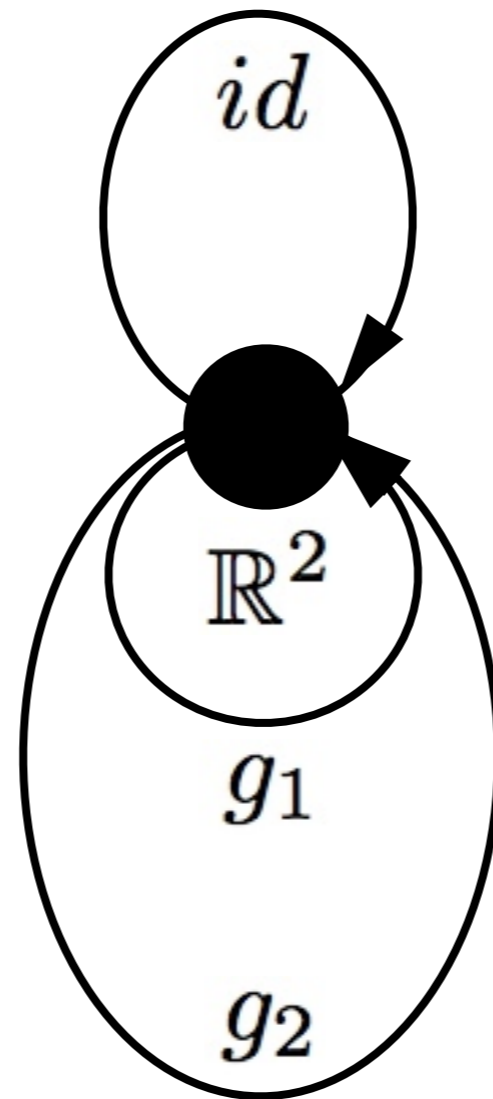
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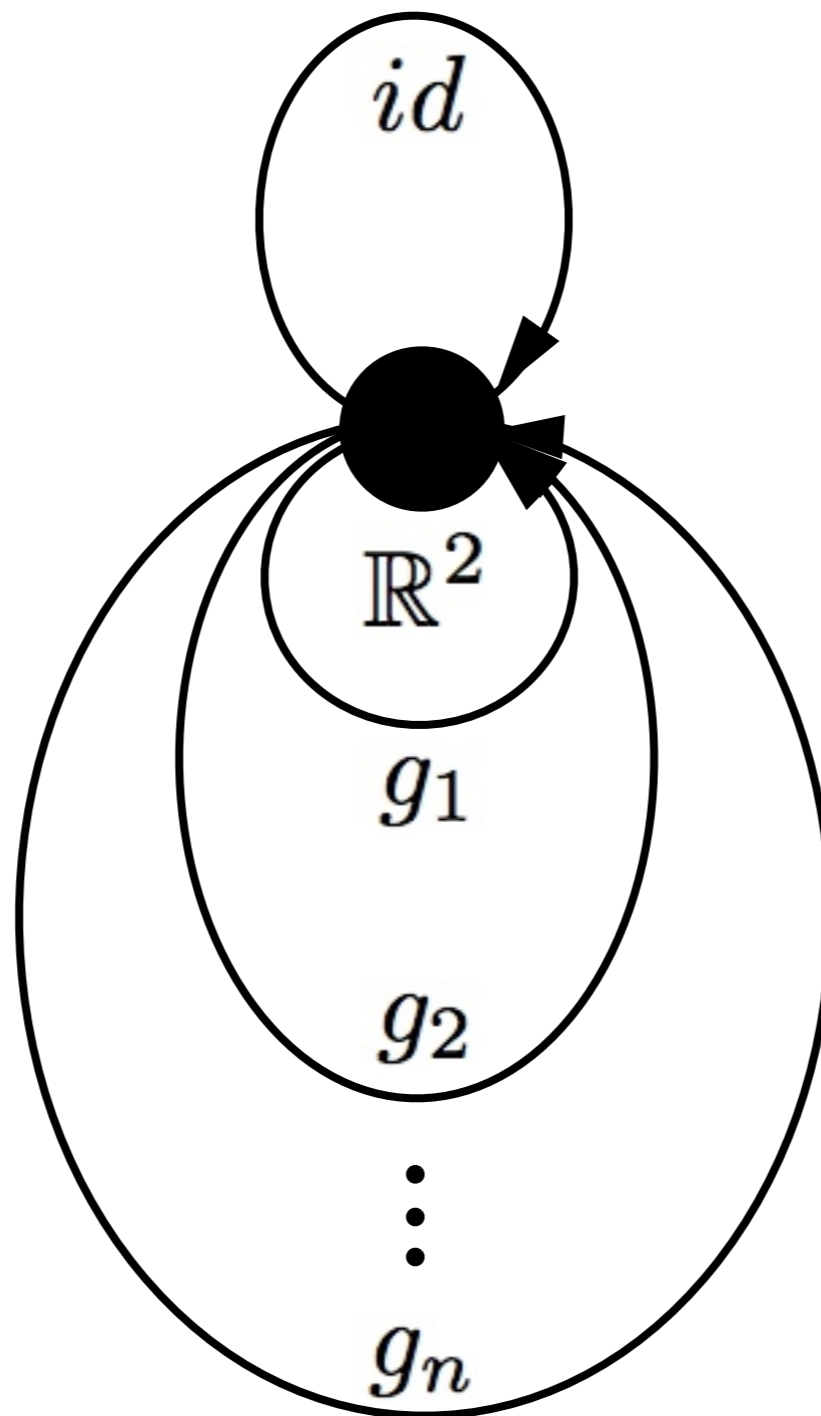
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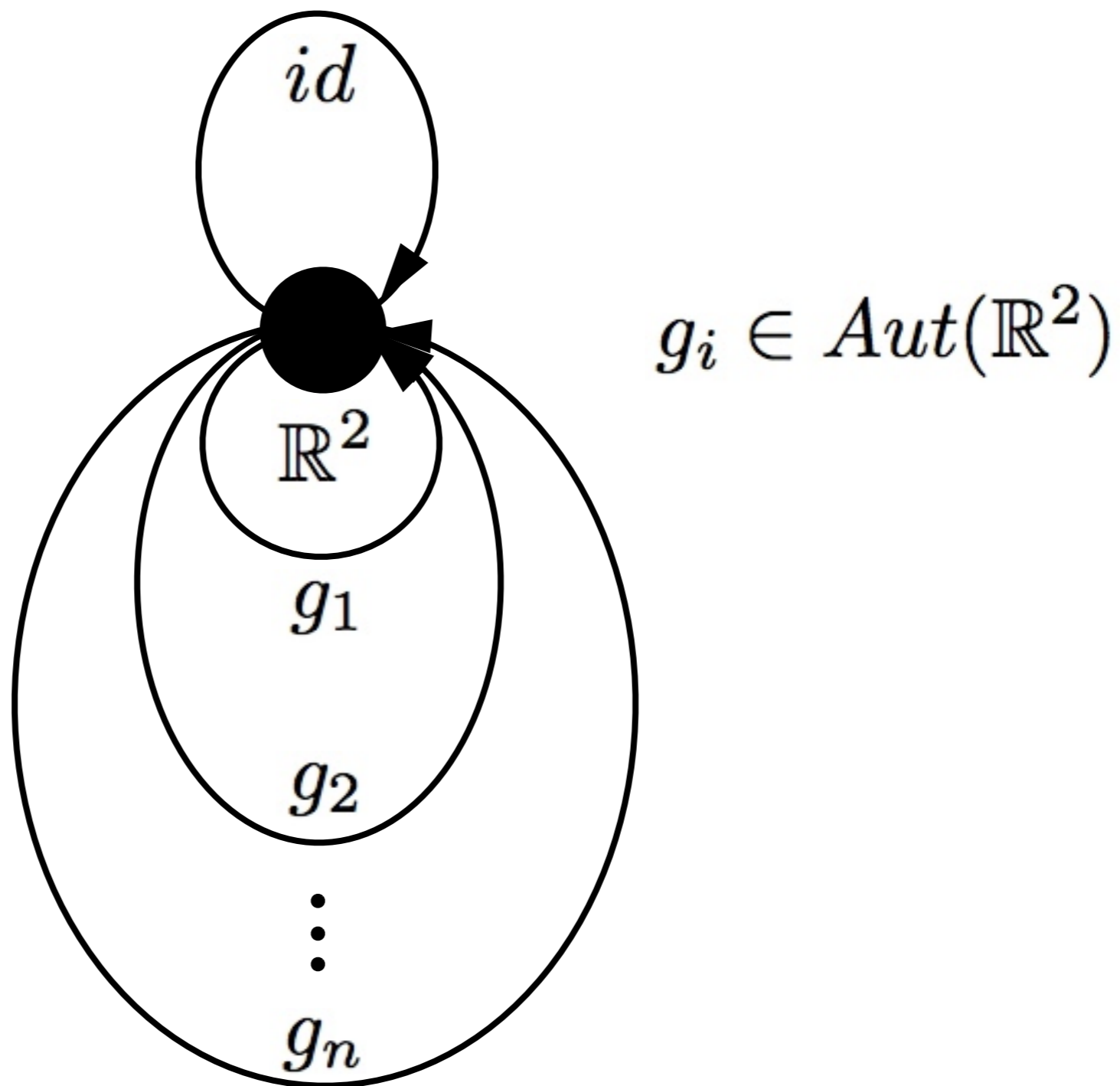
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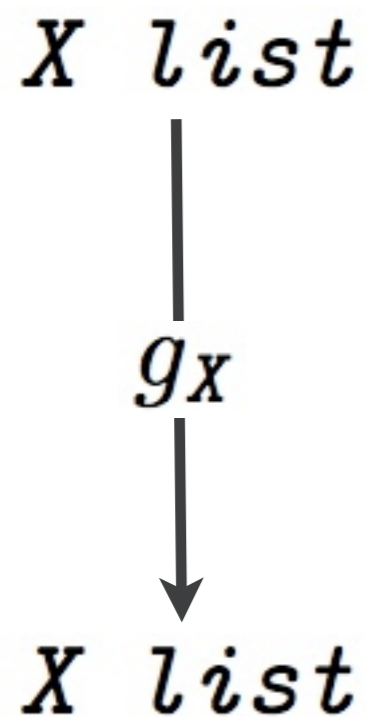
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X list

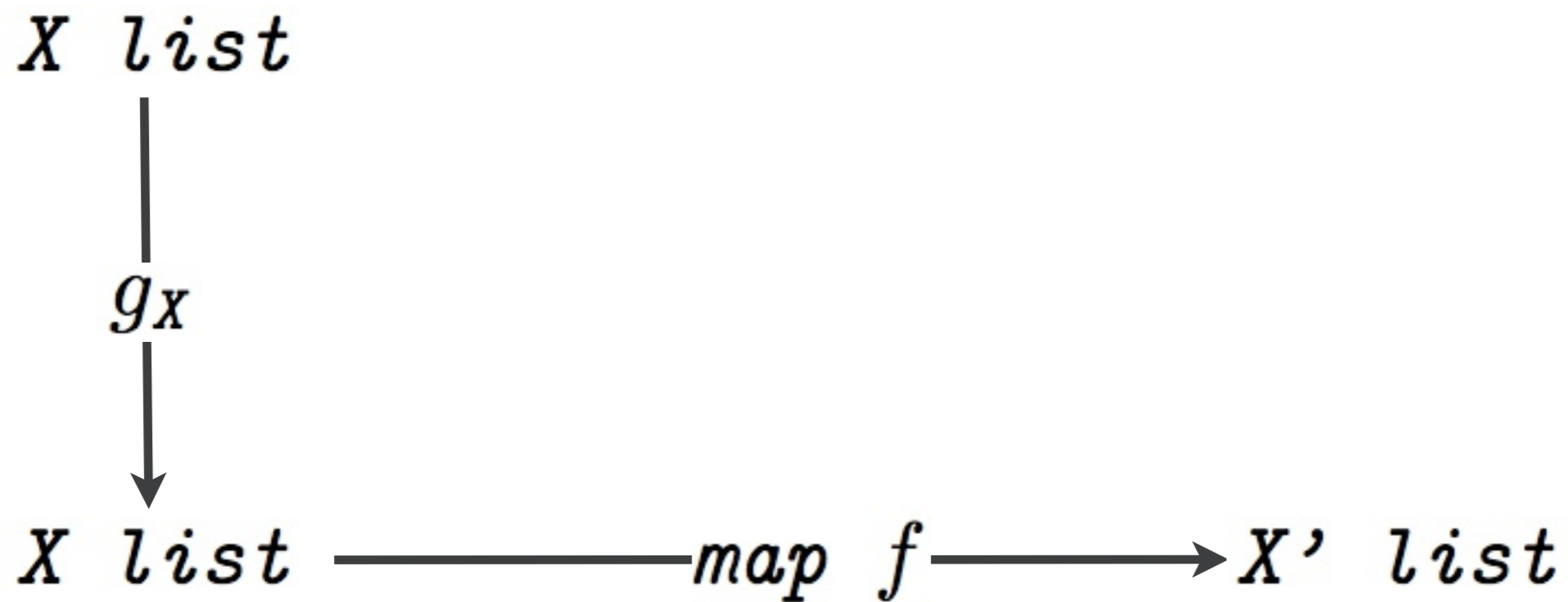
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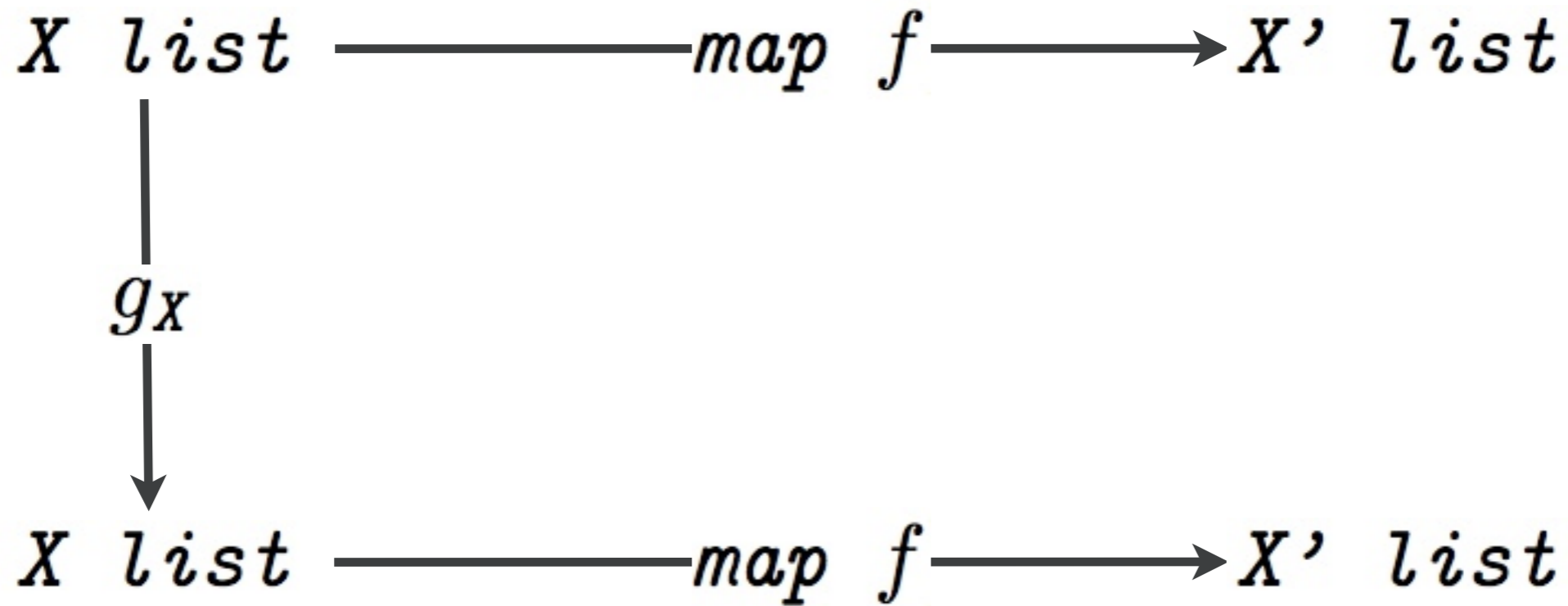
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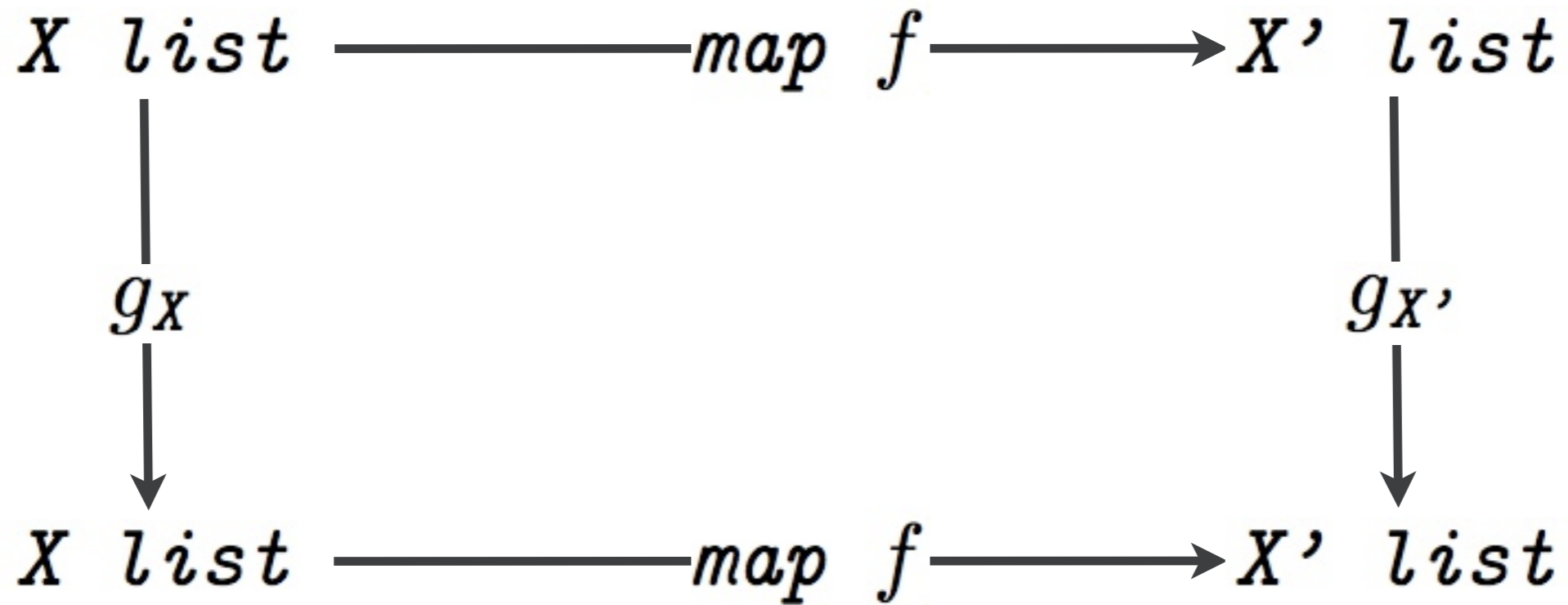
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Atkey: *Main Points*

- Extend System $F\omega$ with type system encoding geometric invariances.
- Interpret kinds as reflexive graphs, types as reflexive graph morphisms.
- Connect free theorems of Wadler/Reynolds with Noether's theorem via symmetries of these reflexive graphs.

Atkey: Takeaways

- Types as geometries is a powerful new way of manipulating our "syntactic discipline".
- Visual intuition, connections to group theory.
- Physics is only one potential application!



∴ τ



∴ τ

The End